

Imaging the Electronic and Vibronic States of Single Semiconductor Nanowires

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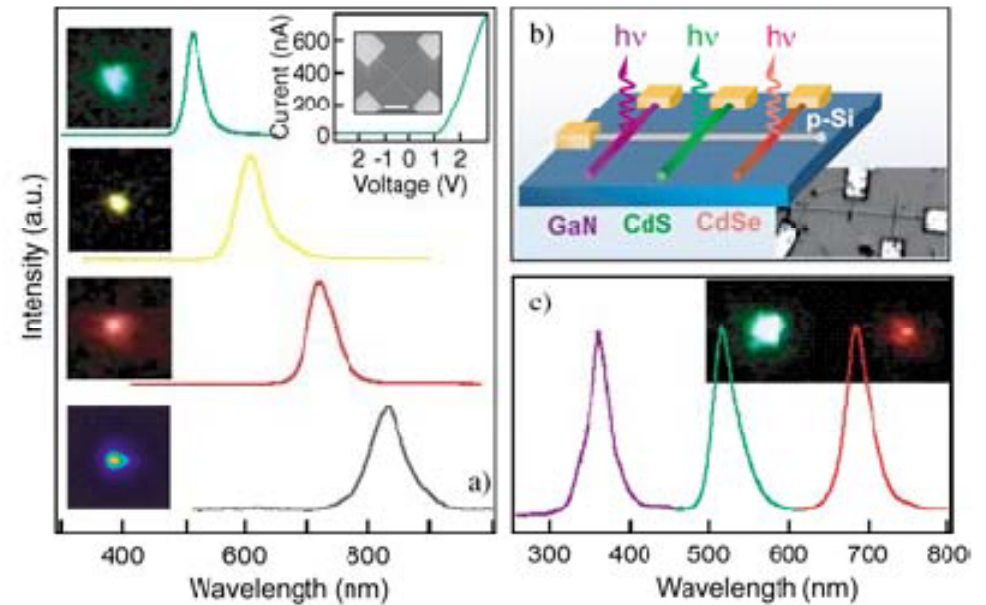
Hannah Joyce

H. Tan

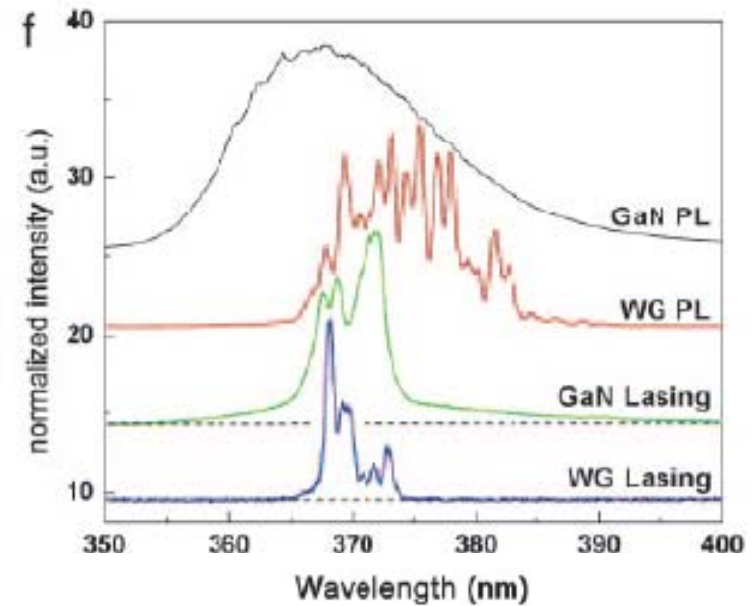
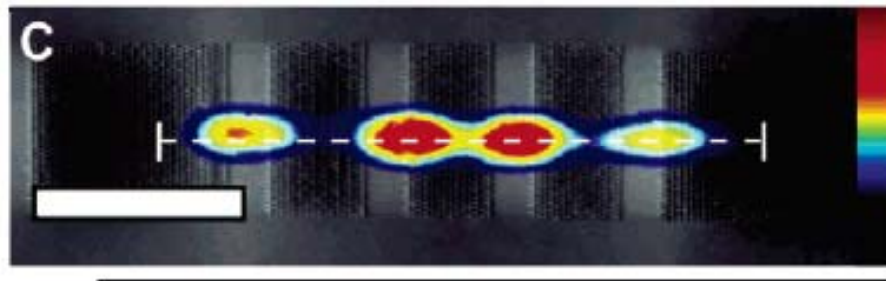
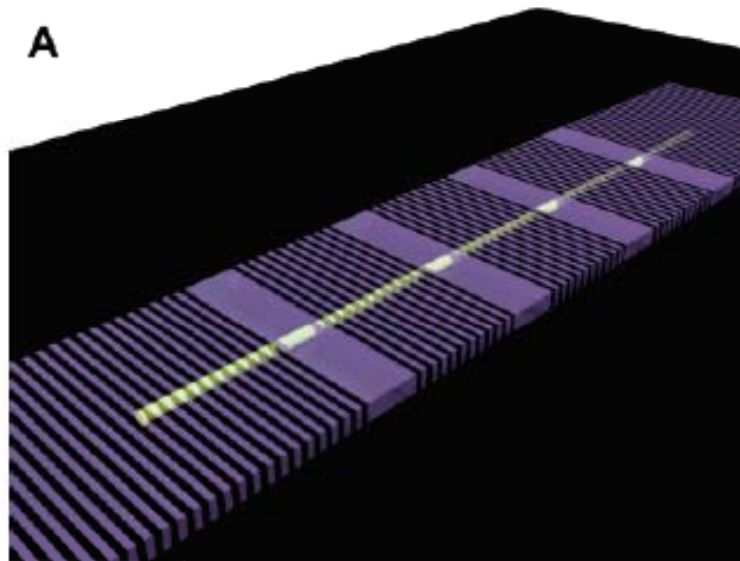
Y. Kim

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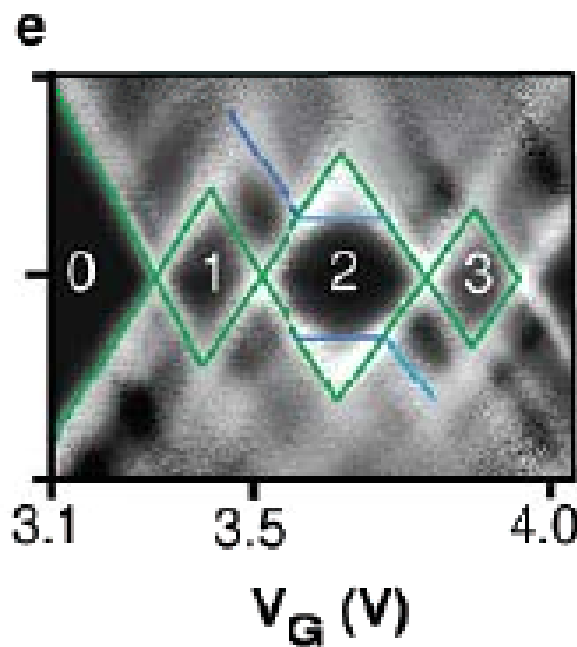
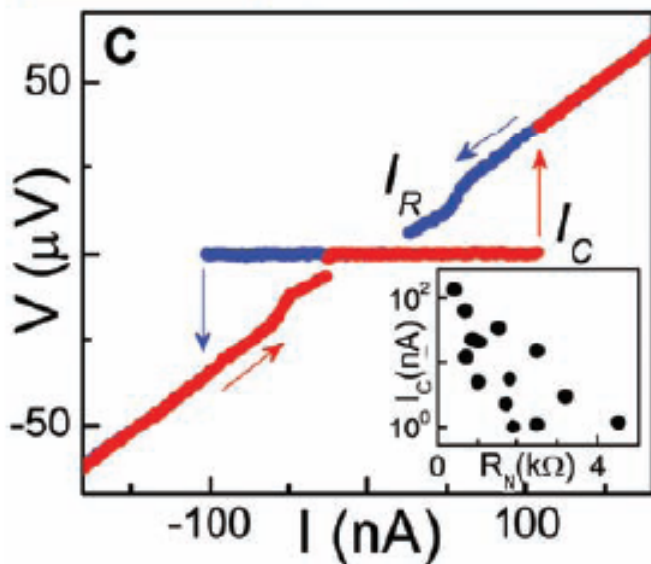
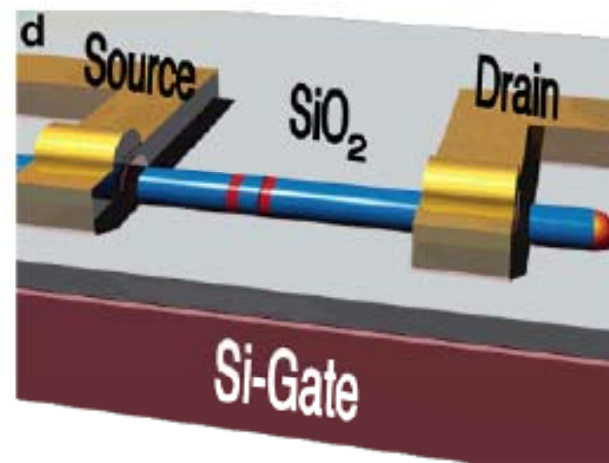
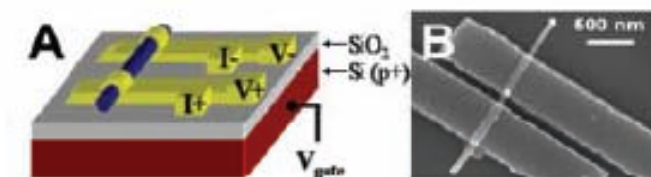
Semiconductor Nanowires as Photodetectors



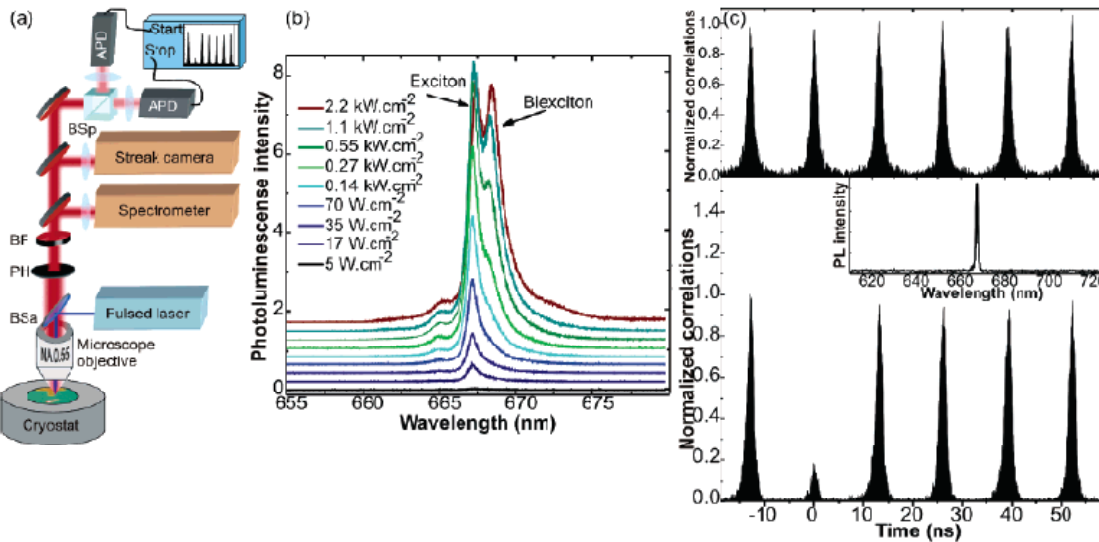
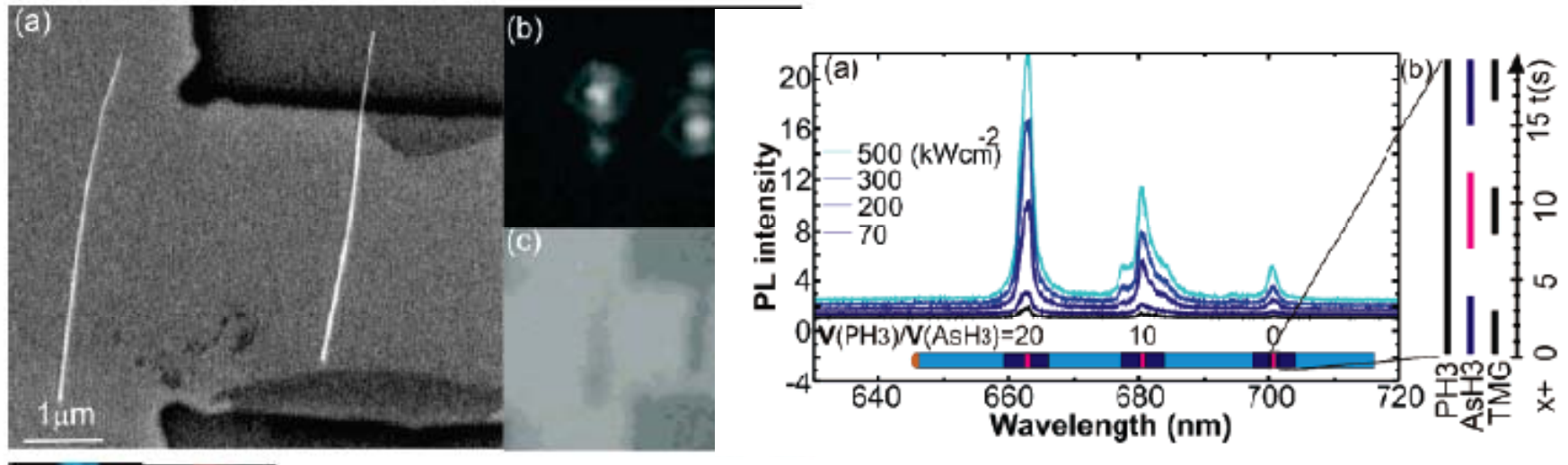
LED and Laser Nanowires



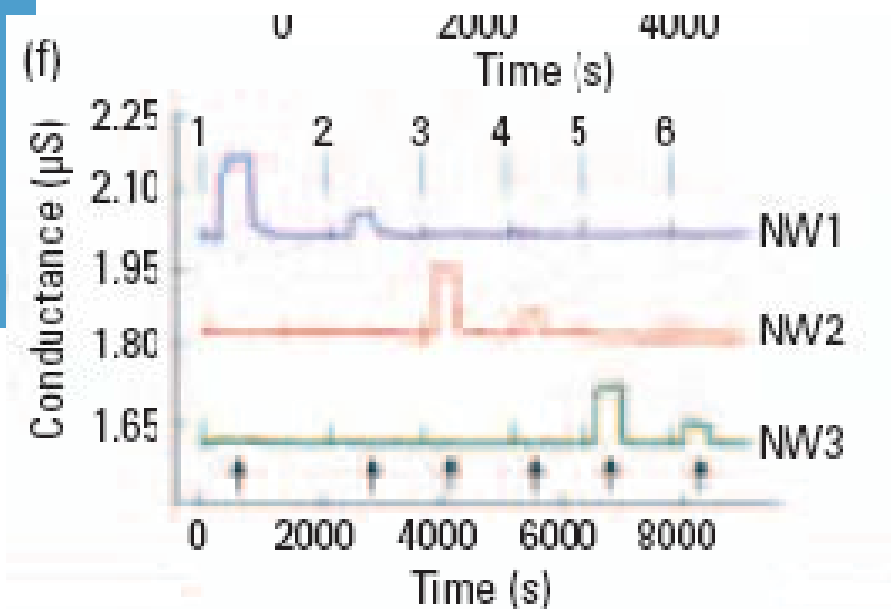
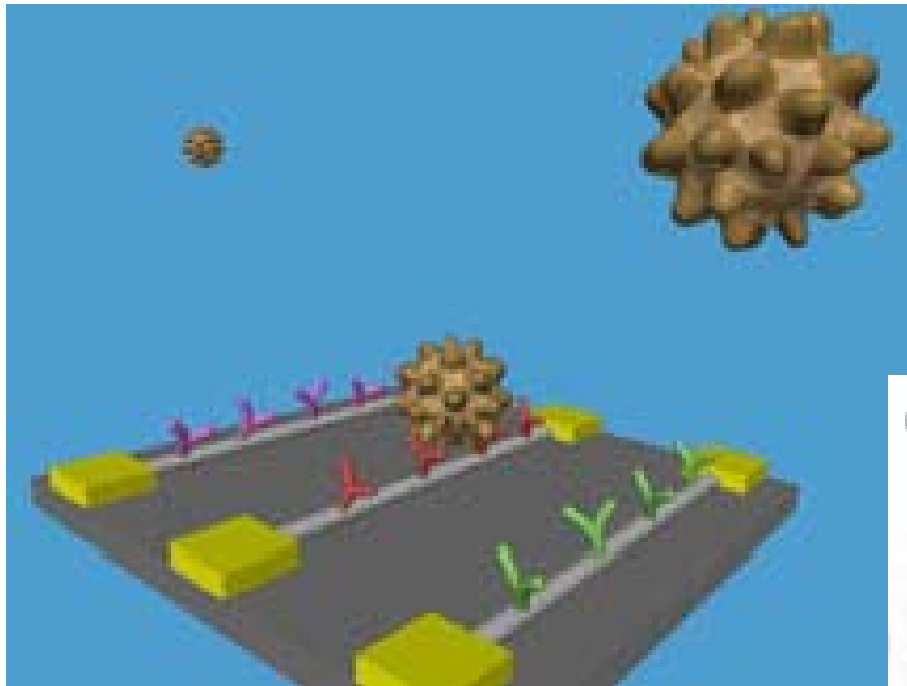
Nanowires as Single Electron Transistors



Nanowires as single photon emitters

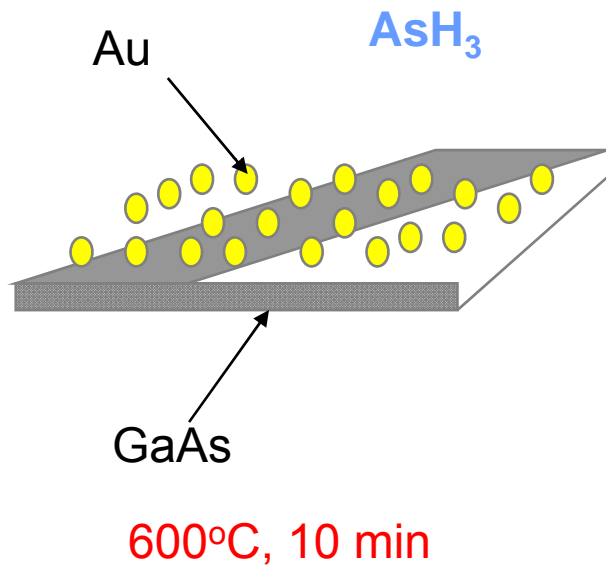


Nanowires as Biosensors



Core-Shell Nanowire Growth

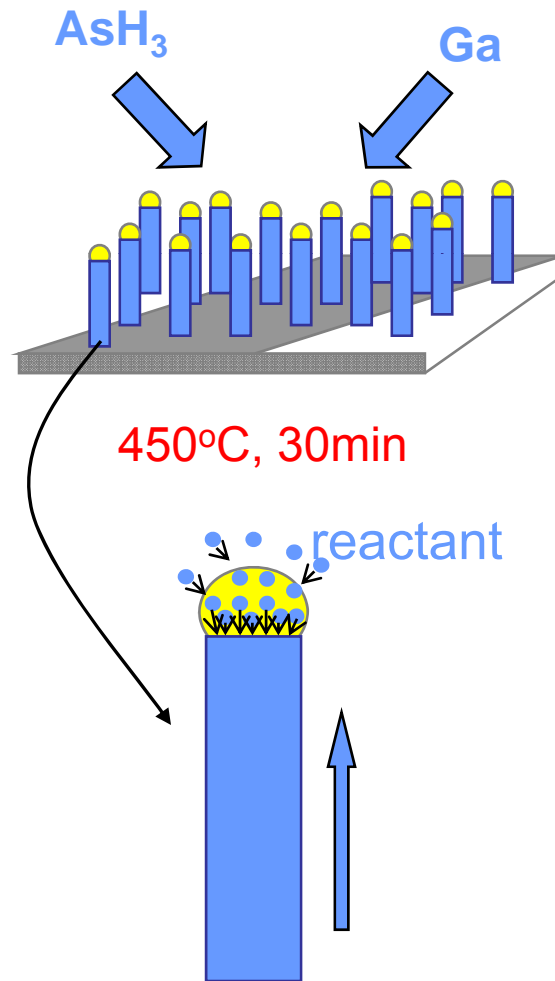
Pre-growth



Desorb surface contaminants and form eutectic alloy.

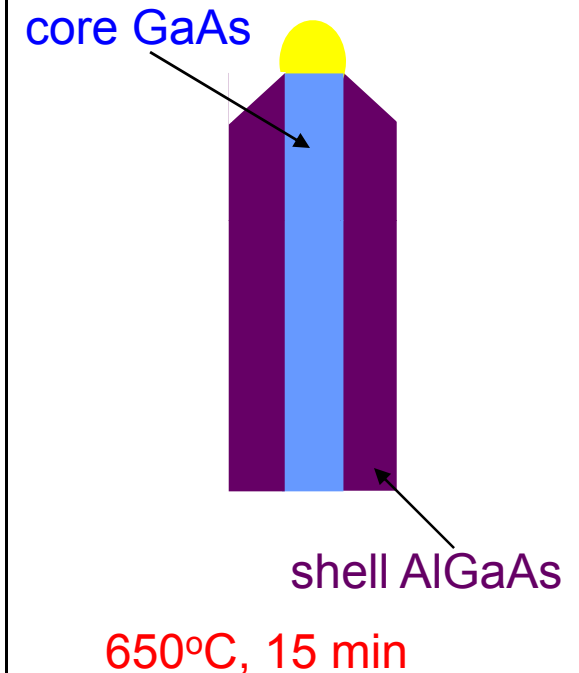
TMS-Orlando, 2007

Core: GaAs



Vapor-Liquid-Solid growth

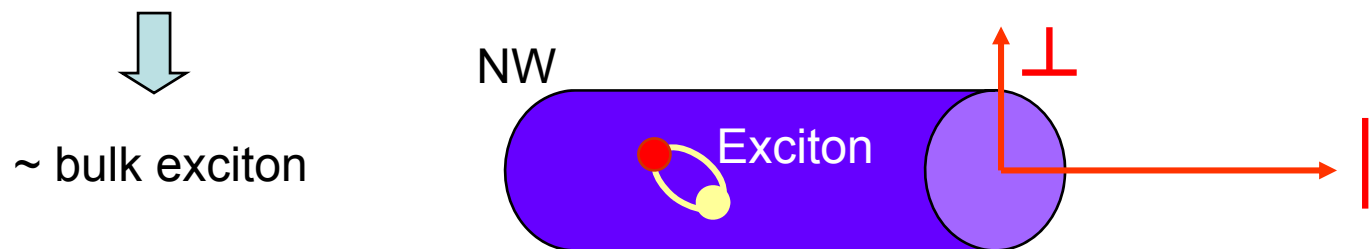
Shell: AlGaAs



Wire diameter is determined by Au catalyst and shell growth time

Motivation

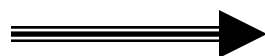
Nanowire diameters D ($\sim 50-150$ nm) $>$ Bohr exciton's diameter (~ 24 nm)



Dielectric "confinement" of EM dipole field ($D \ll \lambda$):

Exciton density

$$N_{\parallel} = N_{\perp}$$



Photoluminescence intensities

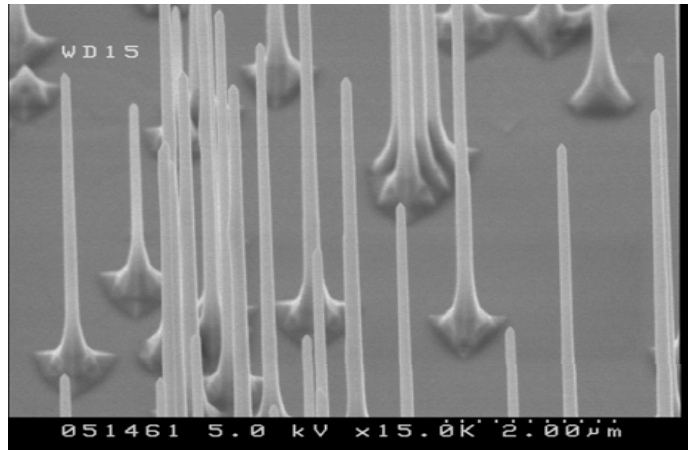
$$I_{\parallel} \gg I_{\perp}$$



$$\tau_{\parallel} \ll \tau_{\perp}$$

We are interested in exciton spin dynamics
of single nanowires

Single nanowire studies

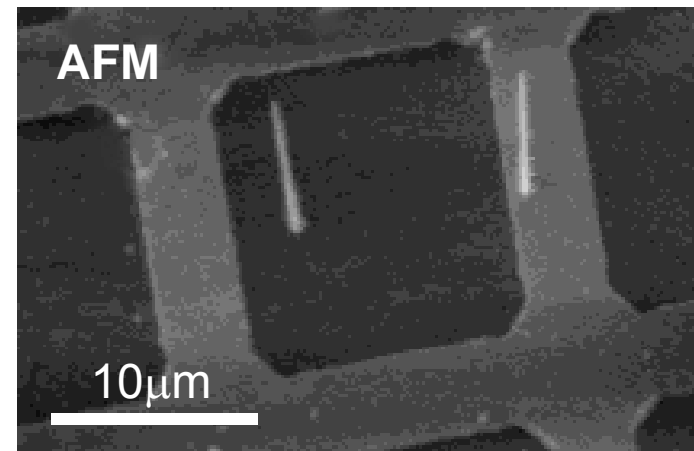


Field-Emission Scanning
Electron Microscope (FESEM)
image

Nanowires were removed from the
growth substrate into solution and
deposited onto a silicon substrate

a single nanowire:

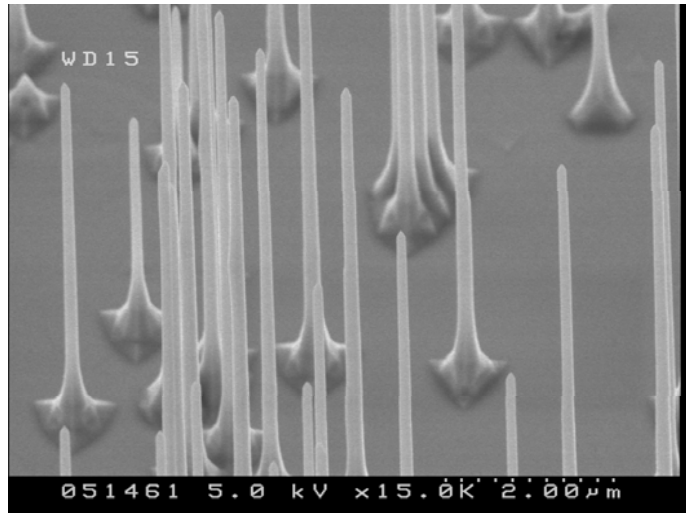
~80nm in diameter, ~5-8 μm long



wire's diameter $>$ Bohr exciton diameter
 \Rightarrow expect no quantum confinement effects



Single nanowire studies

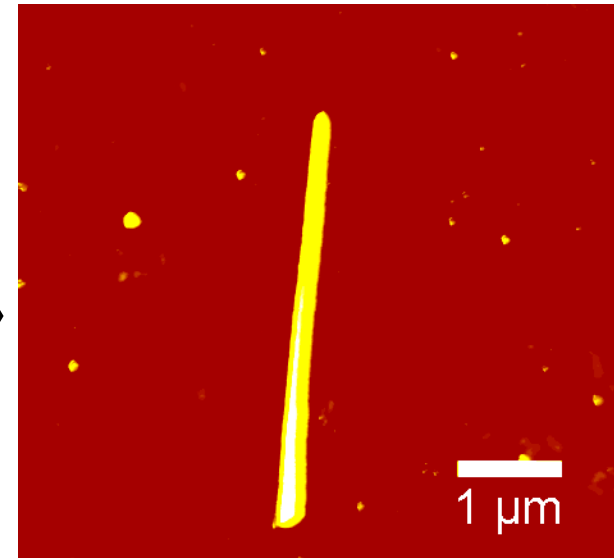


Field-Emission Scanning Electron
Microscope (FESEM) image:
nanowires have tapered shape.

Nanowires were removed from the
growth substrate into solution and
deposited onto a silicon substrate

A single nanowire:

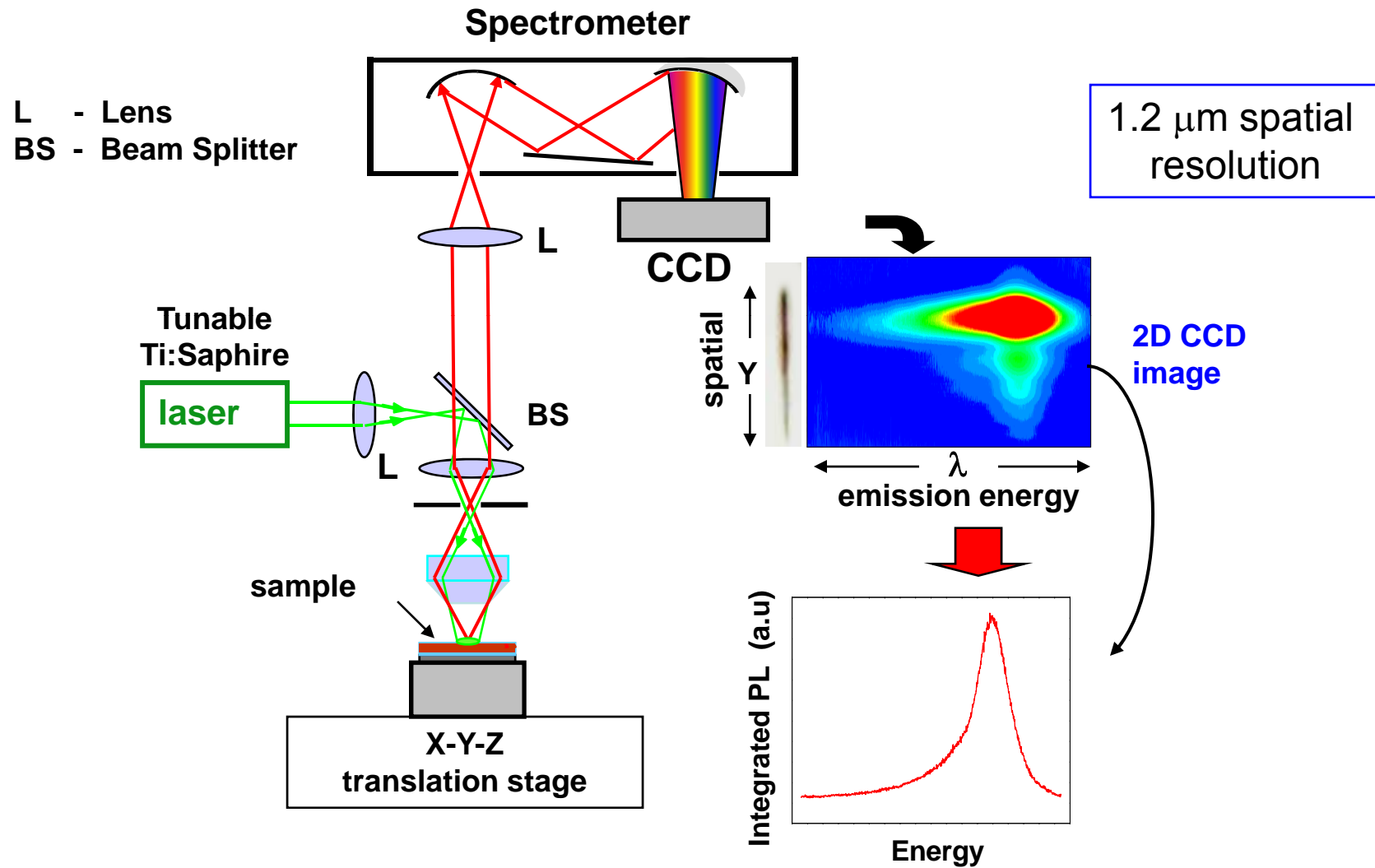
~80nm in diameter, ~5-8 µm long



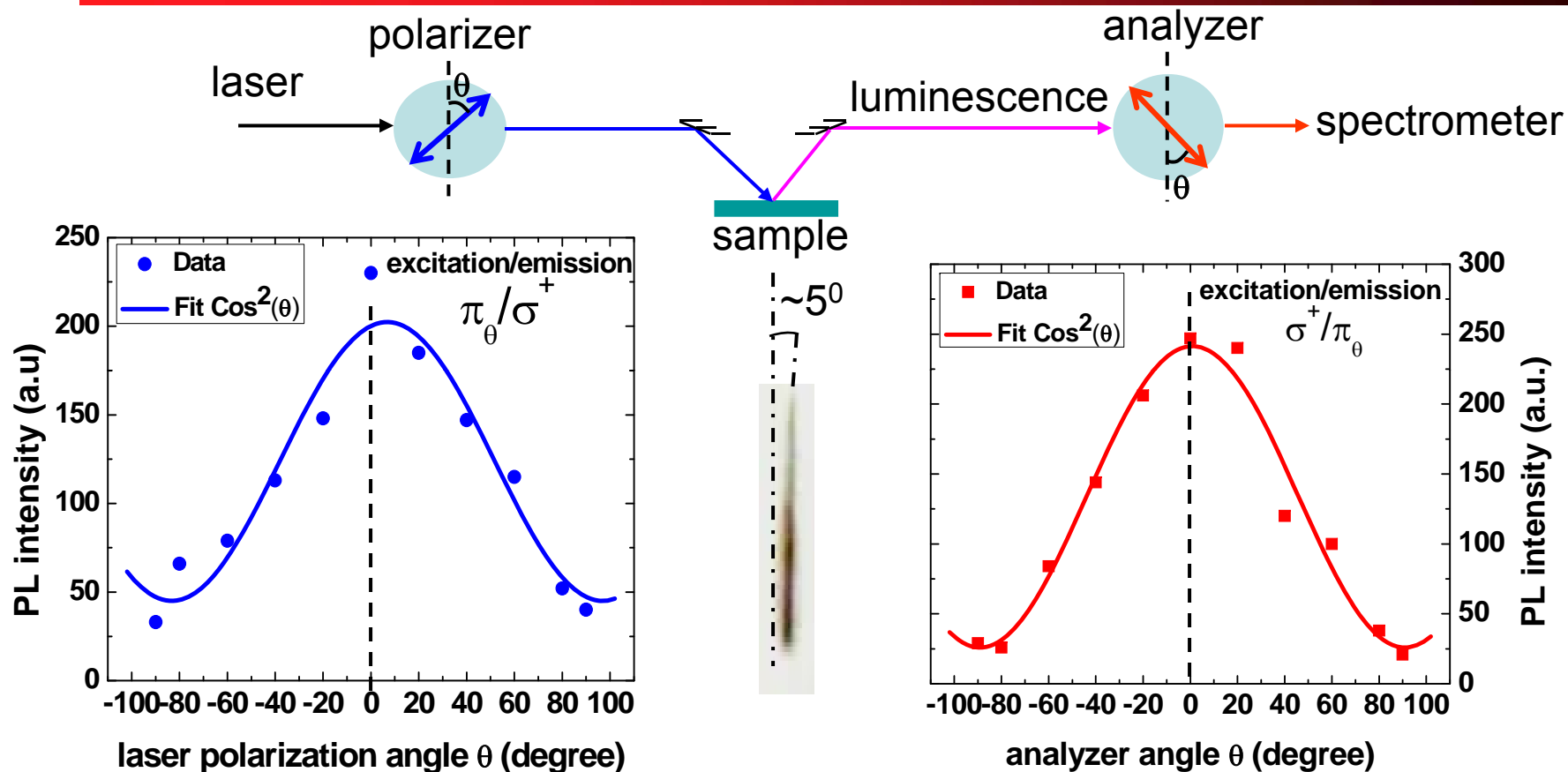
Core diameter > Bohr exciton diameter (24nm)

=> no quantum confinement effects

Photoluminescence Imaging



Polarization studies

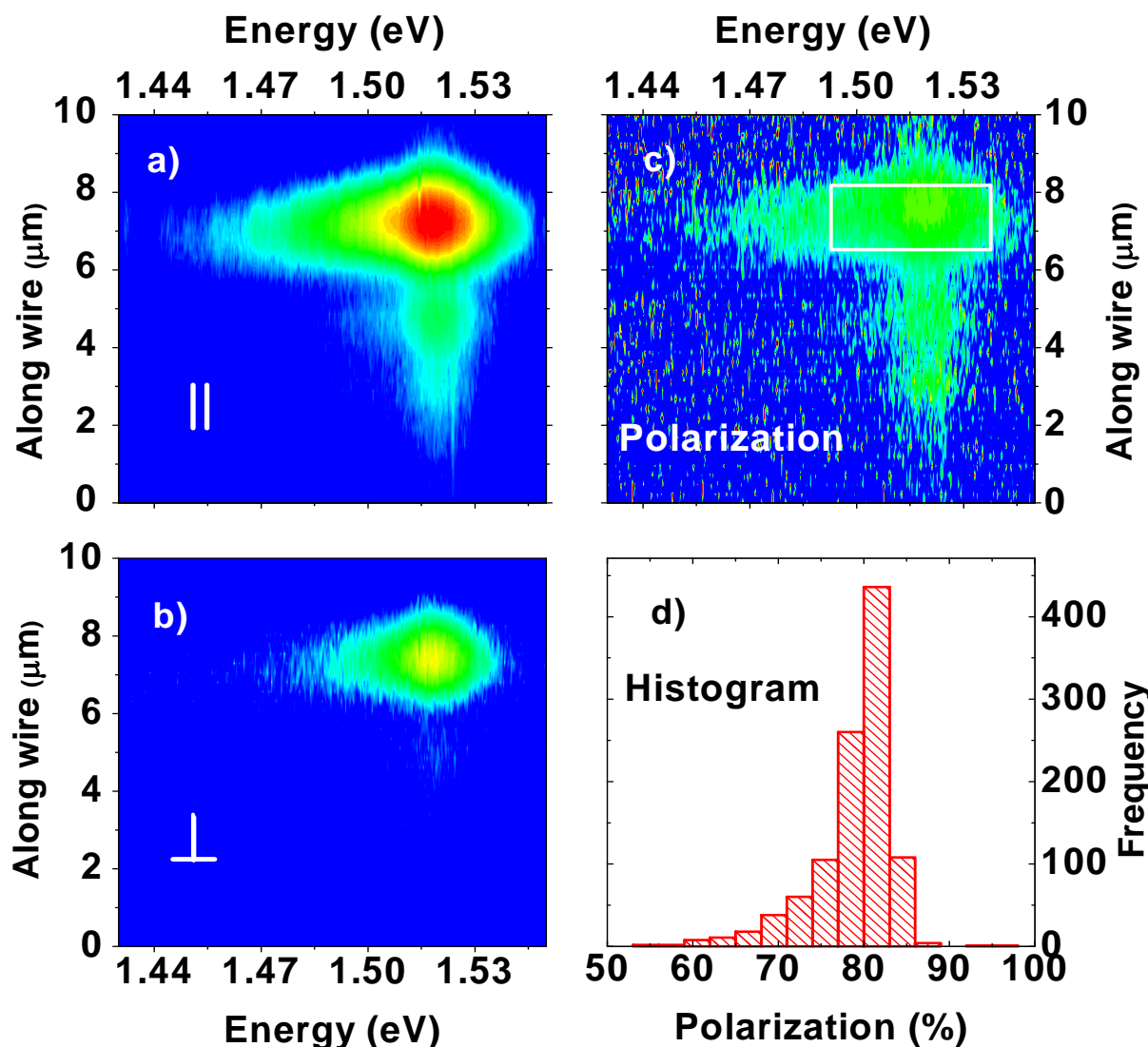


polarizer = π_θ ; analyzer = σ^+

polarizer = σ^+ ; analyzer = π_θ

PL emission is *strongly polarized* parallel to the wire, and is *strongly enhanced* when the laser excitation is polarized parallel to the wire

Polarization Imaging



Calculate pixel by pixel

$$P = \frac{I_{\parallel} - I_{\perp}}{I_{\parallel} + I_{\perp}}$$

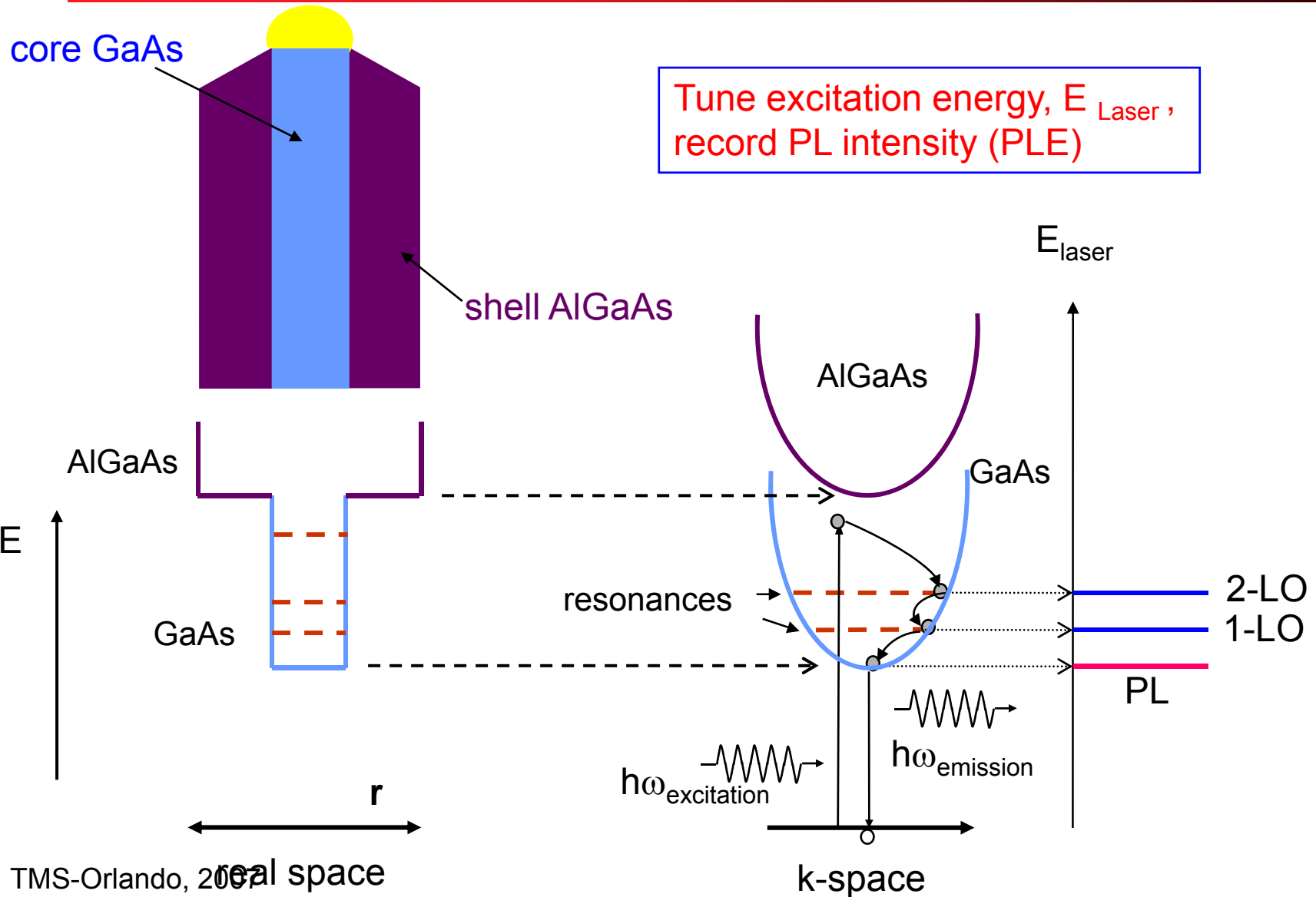
~82%

Strongly polarized due to the large dielectric mismatch between GaAs and air

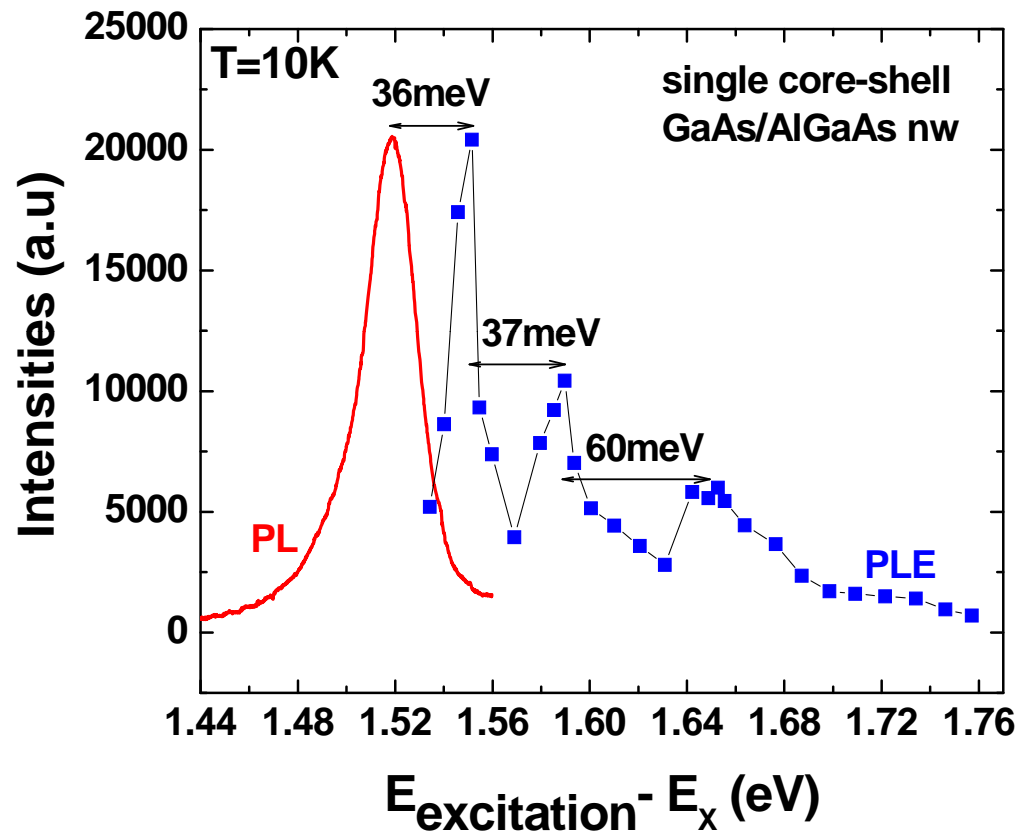
(*Science* 293 1455 (2001),

APL. 89 173126 (2006))

Resonant Excitation

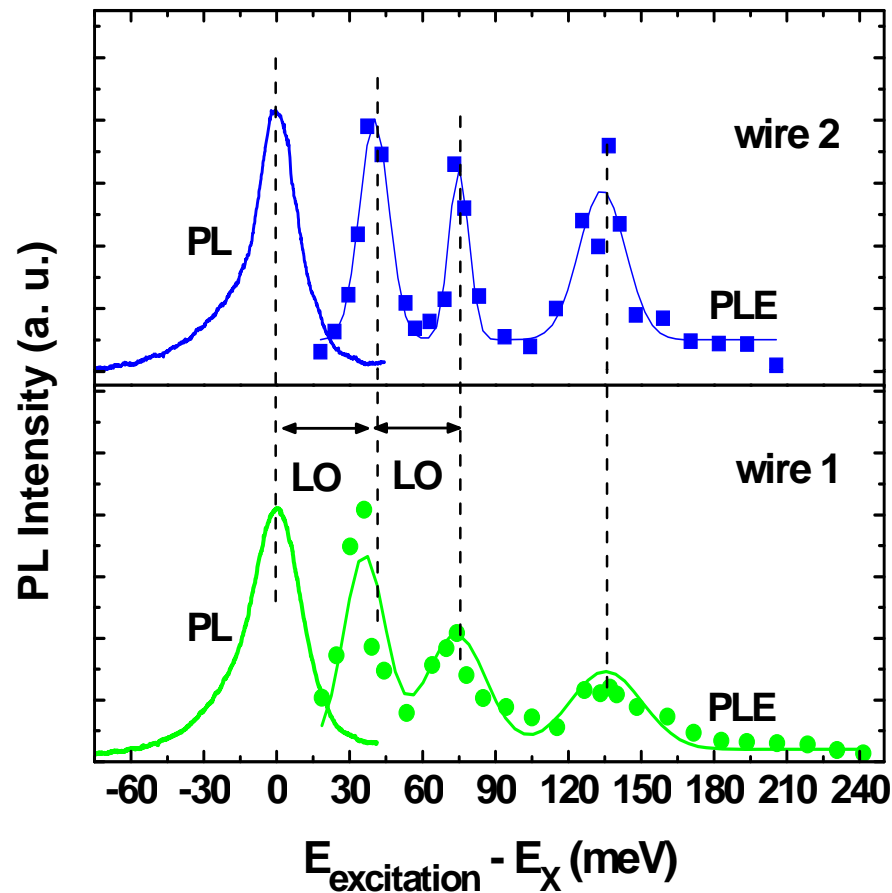


Resonant Excitation



Clear resonances at 36, 73 and ~133 meV above free exciton energy.

Resonant Excitation



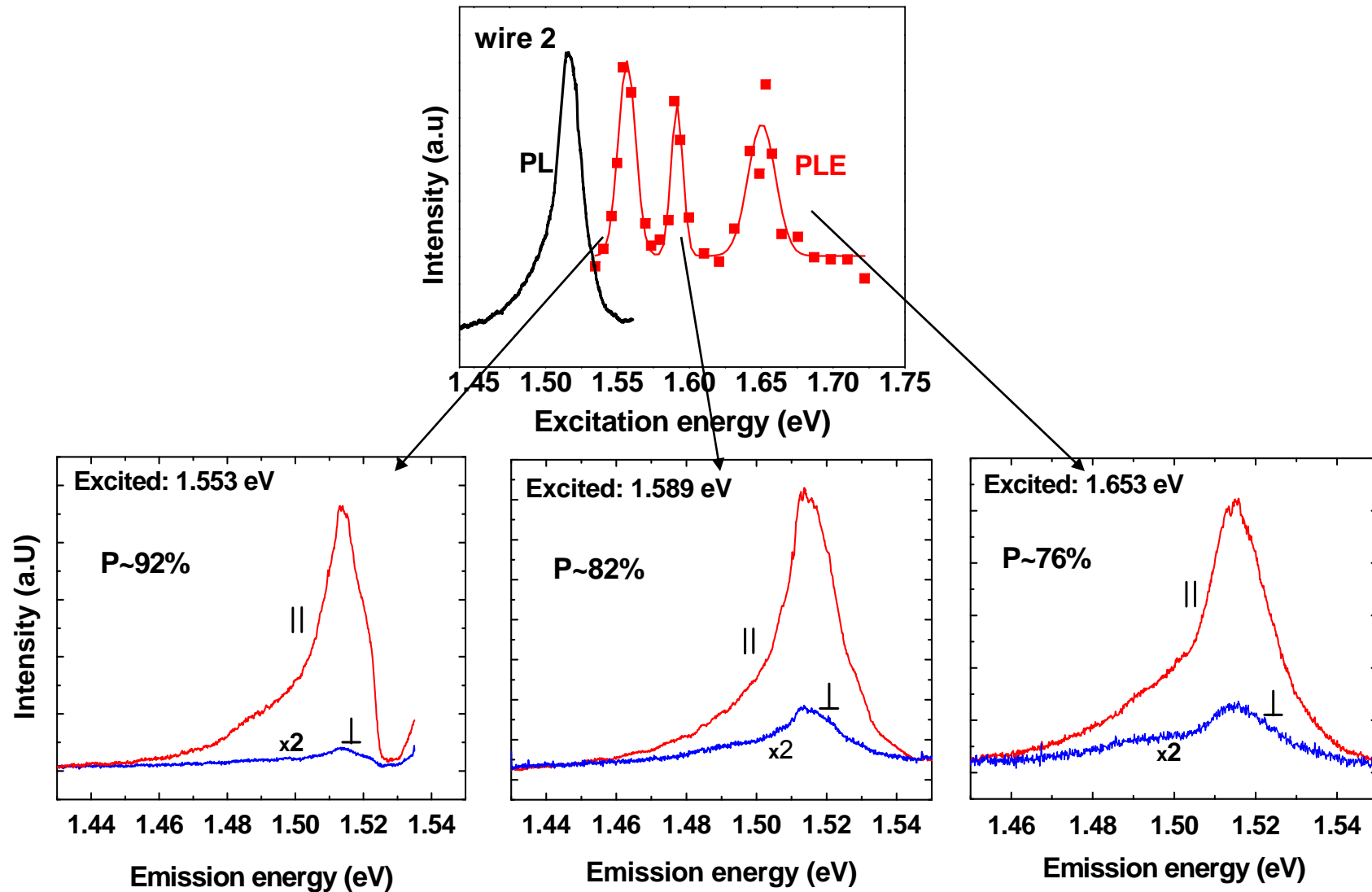
1-LO and 2-LO GaAs phonons

Resonance at ~133 meV:

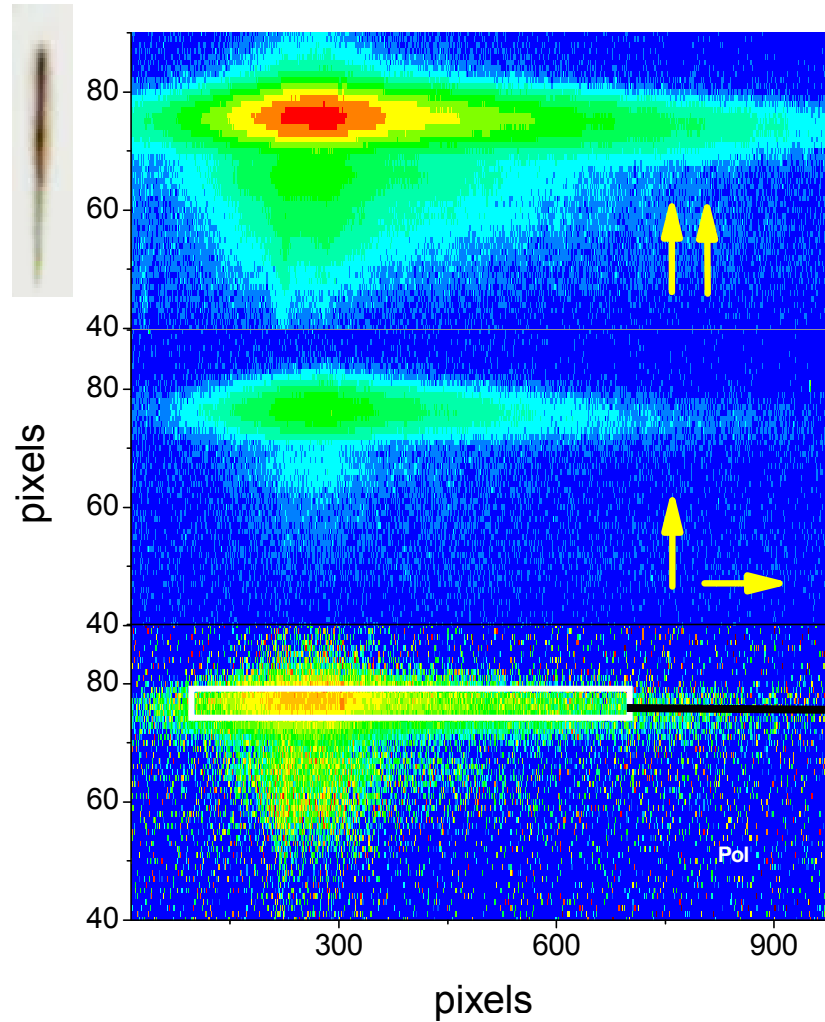
1. Defect-AlGaAs related.
2. Bottom of AlGaAs band (Low concentration of Al ~10%, instead of growth condition 26%)

How does the polarization depend on excitation energy?

Excitation dependent polarization



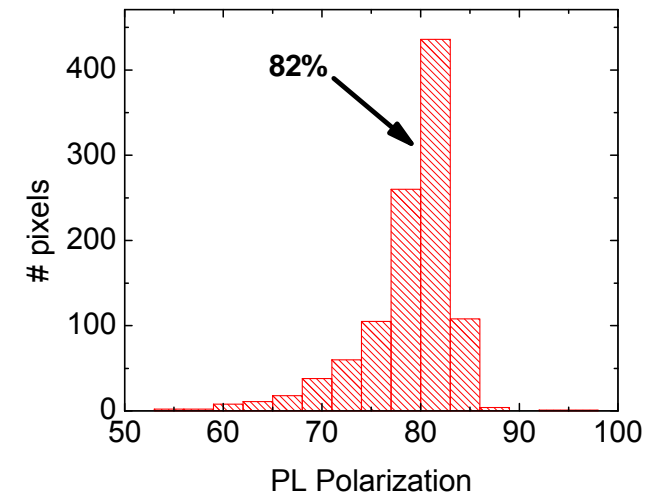
PL Polarization Imaging



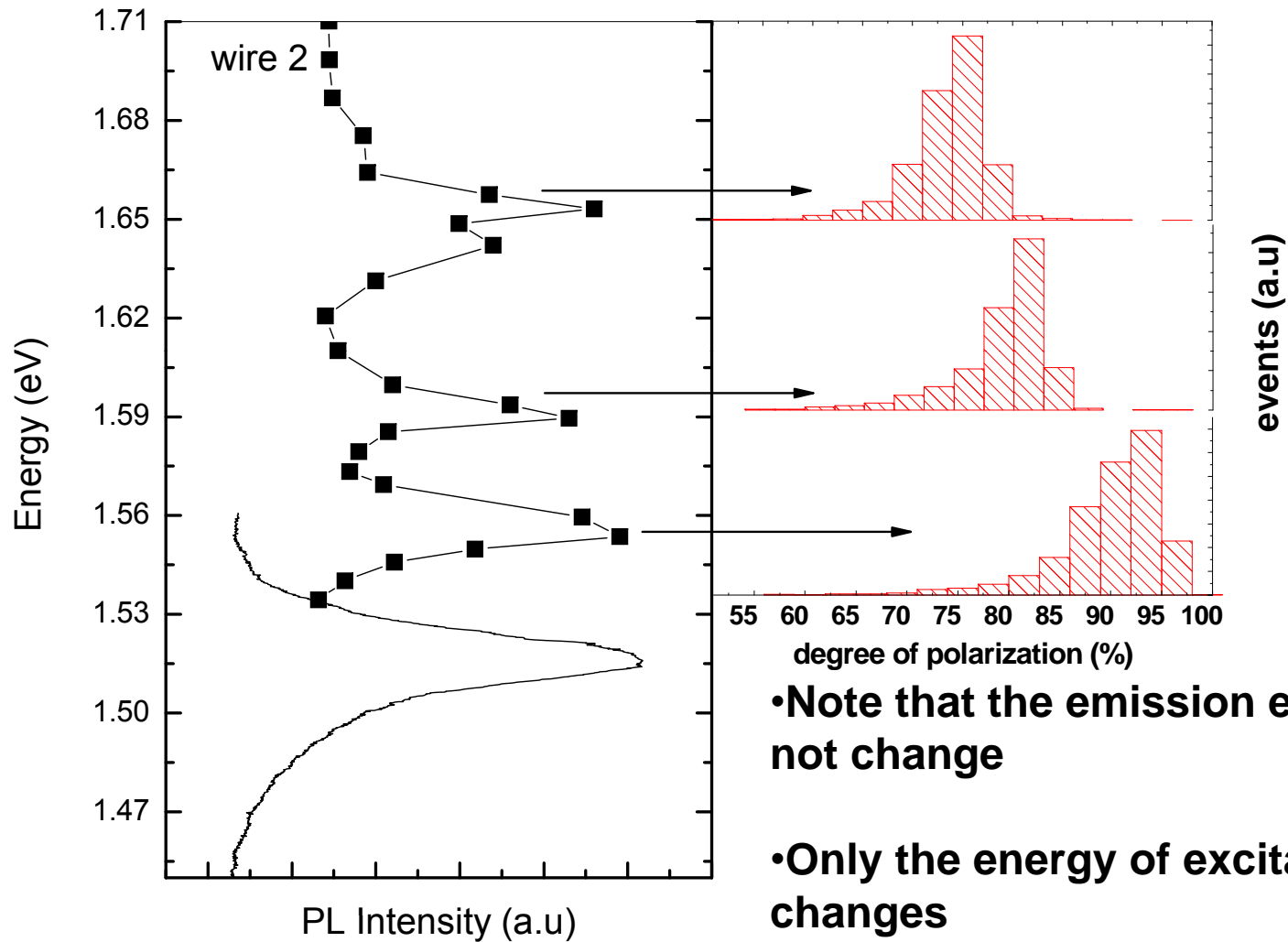
•Excitation laser polarized along nanowire

•Analyze emission polarization

$$P = \frac{I_{\parallel} - I_{\perp}}{I_{\parallel} + I_{\perp}}$$



Polarization depends on excitation energy

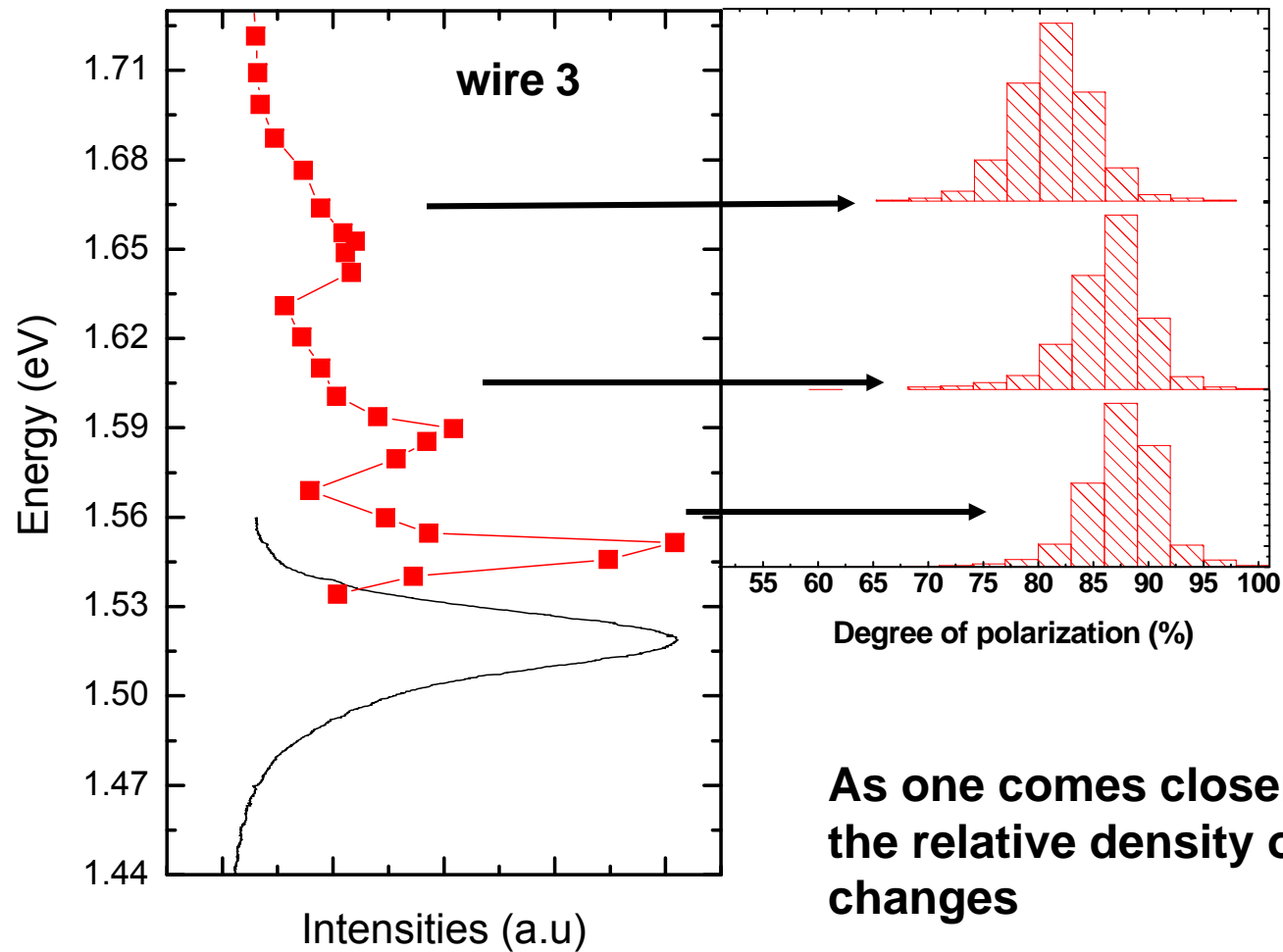


•Note that the emission energy does not change

•Only the energy of excitation changes

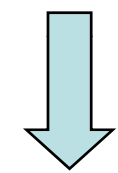
•Changing polarization must result from changing exciton distributions

Polarization excitation dependence also depends on wire...



$$\frac{n_{\parallel}}{n_{\perp}} \rightarrow 1$$

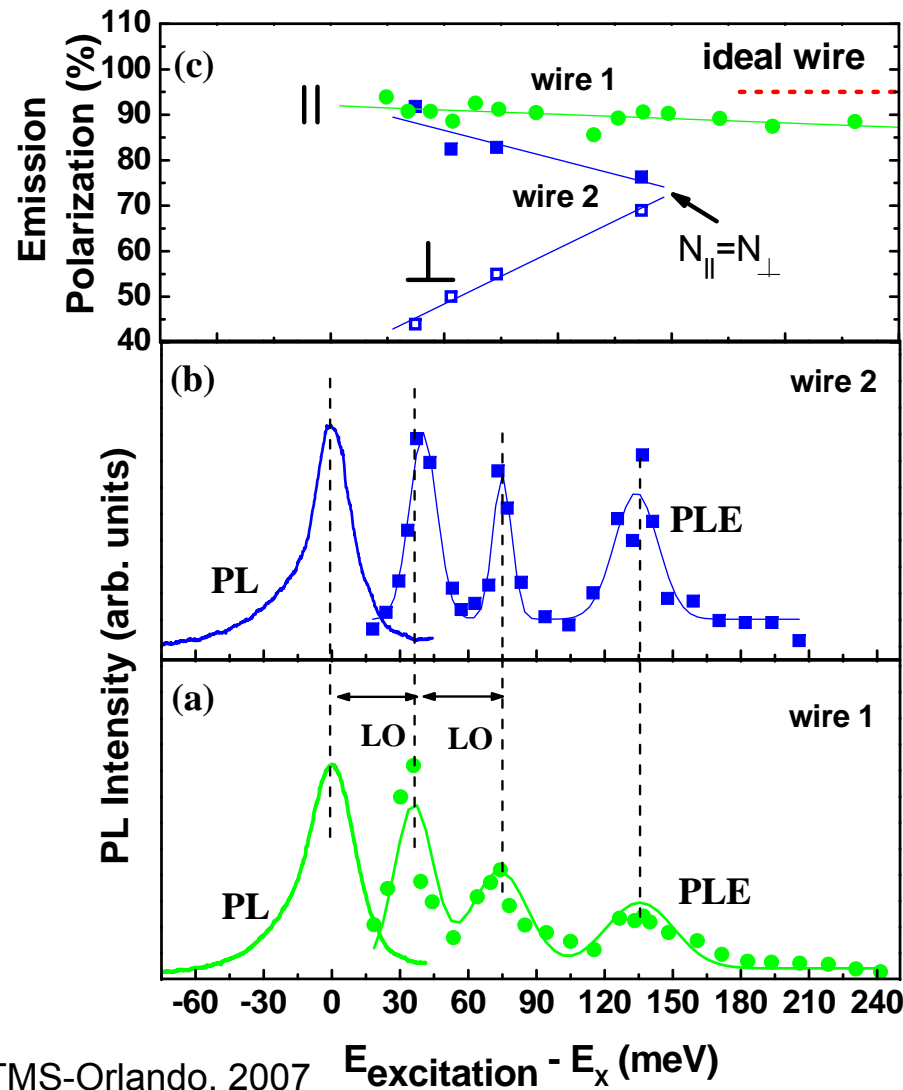
Frequency



$$\frac{n_{\parallel}}{n_{\perp}} > 1$$

As one comes closer to resonance the relative density of excitons changes

Resonant excitation creates non-equilibrium exciton spin distributions



TMS-Orlando, 2007

➤ As excitation comes closer to free exciton energy:

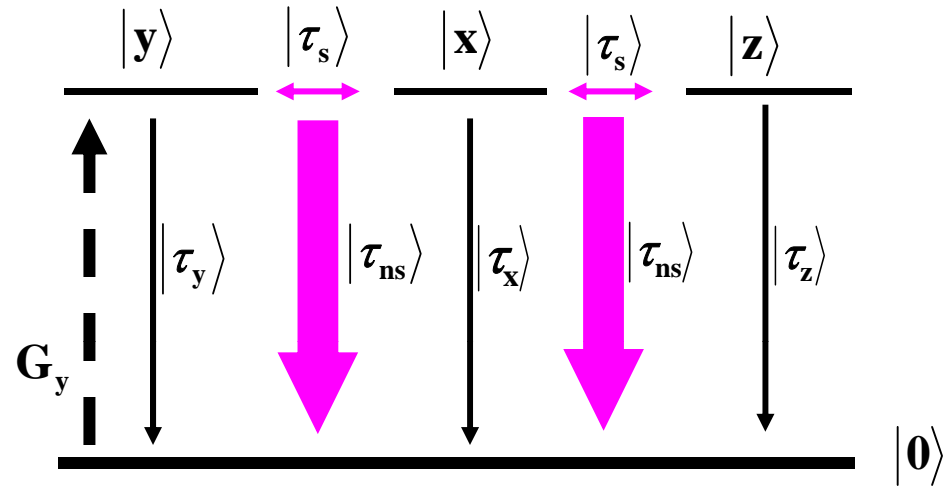
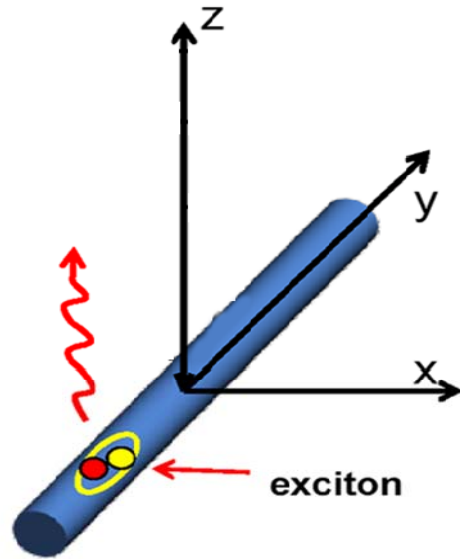
- Along wire: polarization *increases*
- Perpendicular: polarization *decreases*

➤ Polarization are different for different wires

➤ Wire 2: thermal equilibrium

$$N_{\parallel} = N_{\perp}$$

Exciton Dynamics



$$\tau_{x,z} = \tau_y \left(\frac{1 + \epsilon_s}{2} \right)^2$$

$$\tau_y = \tau_{vac} = \frac{3\pi\epsilon_0\hbar c_0^3}{\omega_{exc}^3 D_{exc}^2}$$

$$\tau_{x,z} \gg \tau_y \gg \tau_{nr}, \tau_s \quad \text{and} \quad \frac{I_{\perp}}{I_{\parallel}} \ll 1$$

At thermal equilibrium (highest energies) assume:

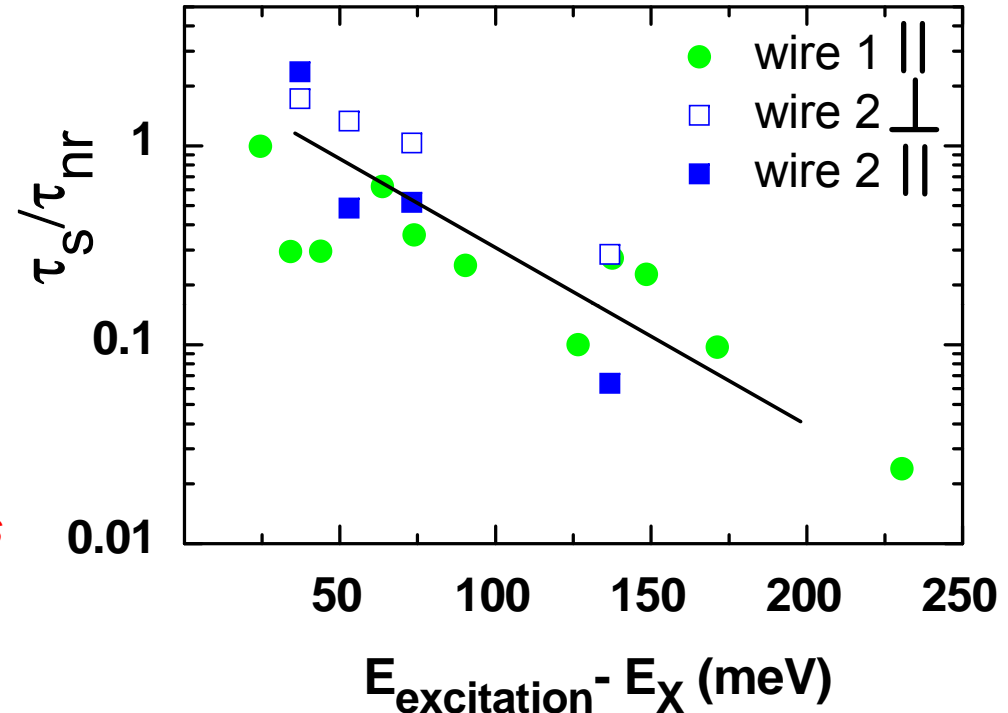
$$n_x = n_y \implies \frac{I_{\perp}}{I_{\parallel}} = \frac{\tau_y}{\tau_x}$$

Spin scattering time

Steady state: $dn_{\alpha} / dt = 0$

$$\frac{\tau_s}{\tau_{nr}} = \frac{I_{\perp}(1+P)}{I_{\parallel}(1-P)} - 1 \quad \text{for } \parallel$$
$$\frac{\tau_s}{\tau_{nr}} = \frac{I_{\perp}(1-P)}{I_{\parallel}(1+P)} - 1 \quad \text{for } \perp$$

**Spin relaxation time depends
on excitation energy**



“Non-Equilibrium Exciton Spin Dynamics in Resonantly Pumped Single Core-Shell GaAs-AlGaAs Nanowires”

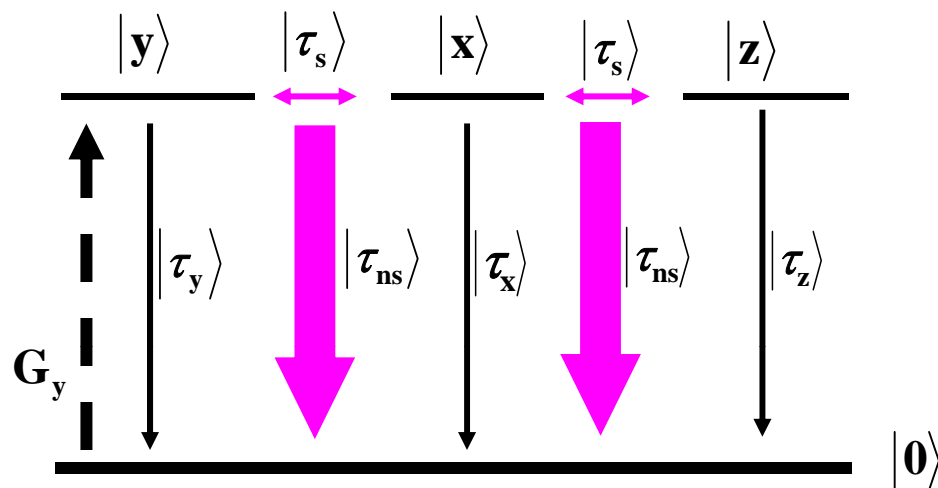
Thang. B. Hoang, L.V. Titova, J. M. Yarrison-Rice, H. E. Jackson, A. O. Govorov, Y. Kim, H. J. Joyce, H. H. Tan, C. Jagadish, L. M. Smith

Conclusions

Single GaAs-AlGaAs NWs under resonant excitation:

- Resonances observed at *1-LO and 2-LO and $\sim 133\text{meV}$* (AlGaAs related) above the PL emission line
- Resonant excitation *creates non-equilibrium exciton dipole distributions*
 - Polarization of PL is strongly enhanced as excitation energy comes closer to resonance with free exciton emission.
 - Rate equations: *dependent of spin relaxation time on excitation energy*

Rate equations



$$\frac{dn_x}{dt} = G_x - \frac{n_x}{\tau_x} - \frac{n_x}{\tau_{nr}} - 2\frac{n_x}{\tau_s} + \frac{n_y}{\tau_s} + \frac{n_z}{\tau_s},$$

$$\frac{dn_y}{dt} = G_y - \frac{n_y}{\tau_y} - \frac{n_y}{\tau_{nr}} - 2\frac{n_y}{\tau_s} + \frac{n_x}{\tau_s} + \frac{n_z}{\tau_s},$$

$$\frac{dn_z}{dt} = -\frac{n_y}{\tau_z} - \frac{n_y}{\tau_{nr}} - 2\frac{n_y}{\tau_s} + \frac{n_x}{\tau_s} + \frac{n_z}{\tau_s},$$