Consider eq. (13.31). On the l.h.s. m/d ordinarily goes to ρ . However, if we have a concentrated mass, we, instead, cancel d in m/d on the l.h.s. and τ/d on the r.h.s. Then we take the limit of $d \to 0$ and, as a result, obtain the discontinuity in the derivative, namely

$$\frac{\partial q}{\partial x}\mid_{0+} -\frac{\partial q}{\partial x}\mid_{0-} = \frac{m}{\tau} \ddot{q}\mid_{0}$$

On the other hand, per(13.33),

$$\ddot{q} = \frac{\rho}{\tau} \frac{\partial^2 q}{\partial x^2} = v^2 \frac{\partial^2 q}{\partial x^2}$$

Going to the notations of §13.7, we find solution the wave equation

$$\frac{\partial^2 \psi}{\partial x^2} + k^2 \psi = 0, \ k^2 = \frac{\omega^2}{v^2}$$

subject to the boundary conditions

$$\psi_2(0) = \psi_1(0)$$

$$\frac{\partial \psi_2}{\partial x} \mid x=0 - \frac{\partial \psi_1}{\partial x} \mid_{x=0} = -\frac{M}{\tau} v^2 k^2 \psi(0) = -\frac{M\omega^2}{\tau} \psi(0)$$

(notice that per first b.c. it does not matter whether we use $\psi_2(0)$ or $\psi_1(0)$ on the r.h.s. of the second), where

$$\psi_1(x) = \exp\left(ikx\right) + B\exp\left(-ikx\right)$$

is the incident and reflected waves and

$$\psi_2\left(x\right) = C\exp\left(ikx\right)$$

is the outgoing wave. The reflection and transmission coefficients are $R = |B|^2$ and $T = |C|^2$ respectively.