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we, instead, cancel d in m/d on the l.h.s. and τ/d on the r.h.s. Then we take the limit of $d \to 0$ and, as Consider eq. (13.31). On the l.h.s. m/d ordinarily goes to ρ . However, if we have a concentrated mass, a result, obtain the discontinuity in the derivative, namely

$$
\frac{\partial q}{\partial x}\mid_{0+} -\frac{\partial q}{\partial x}\mid_{0-} =\frac{m}{\tau}\ddot{q}\mid_{0}
$$

On the other hand, per(13.33),

$$
\ddot{q} = \frac{\rho}{\tau} \frac{\partial^2 q}{\partial x^2} = v^2 \frac{\partial^2 q}{\partial x^2}
$$

Going to the notations of $\S 13.7$, we find solution the wave equation

$$
\frac{\partial^2 \psi}{\partial x^2} + k^2 \psi = 0, \ k^2 = \frac{\omega^2}{v^2}
$$

subject to the boundary conditions

$$
\psi_2(0) = \psi_1(0)
$$

\n
$$
\frac{\partial \psi_2}{\partial x} \quad | \quad x = 0 - \frac{\partial \psi_1}{\partial x} |_{x=0} = -\frac{M}{\tau} v^2 k^2 \psi(0) = -\frac{M \omega^2}{\tau} \psi(0)
$$

(notice that per first b.c. it does not matter whether we use $\psi_2(0)$ or $\psi_1(0)$ on the r.h.s. of the second), where

$$
\psi_1(x) = \exp(ikx) + B \exp(-ikx)
$$

is the incident and reflected waves and

$$
\psi_2(x) = C \exp(i k x)
$$

is the outgoing wave. The reflection and transmission coefficients are $R = |B|^2$ and $T = |C|^2$ respectively.