

# Integration

## Math Physics Quiz 2

10-30-2007

1. The concentration of gas in a cylinder of radius  $R$  and height  $h$ , rotating with angular velocity  $\omega$ , is

$$n(r) = C \exp\left(\frac{m\omega^2 r^2}{2T}\right)$$

where  $r$  is the distance from the cylinder's axis,  $m$  is a molecules mass and  $T$  is temperature (in energy units). Given the total number  $N$  of molecules in the cylinder, evaluate the constant  $C$  and find  $n(r)$ . Express your answer in terms of

$$n_0 \equiv \frac{N}{V}$$

where  $V$  is the volume of the cylinder. Defining a dimensionless parameter

$$\alpha \equiv \frac{m\omega^2 R^2}{2T}$$

investigate your answer in the limit  $\alpha \ll 1$  and  $\alpha \gg 1$ .

*Solution*

Normalization gives

$$\begin{aligned} N &= \int_V n(r) dV = 2\pi h C \int_0^R \exp\left(\frac{m\omega^2 r^2}{2T}\right) r dr = 2\pi h C \int_0^{R^2/2} \exp\left(\frac{m\omega^2 z}{T}\right) dz \\ &= \frac{2\pi h C T}{m\omega^2} \left[ \exp\left(\frac{m\omega^2 R^2}{2T}\right) - 1 \right] \end{aligned}$$

and

$$n(r) = \frac{Nm\omega^2}{2\pi hT} \frac{\exp(m\omega^2 r^2/2T)}{\exp(m\omega^2 R^2/2T) - 1} = n_0 \frac{\alpha}{\exp \alpha - 1} \exp\left(\frac{\alpha r^2}{R^2}\right)$$

with the limits

$$\begin{aligned} n(r) &\stackrel{\alpha \ll 1}{\longrightarrow} n_0 \\ n(r) &\stackrel{\alpha \gg 1}{\longrightarrow} n_0 \alpha \exp\left[-\alpha \left(1 - \frac{r^2}{R^2}\right)\right] \end{aligned}$$

2. The partition function of an ideal monatomic gas, confined to a volume  $V$  at temperature  $T$ , is given by

$$Z = \frac{V}{(2\pi\hbar)^3} \int d^3p \exp\left(-\frac{p^2}{2mT}\right)$$

where  $m$  is the mass of each atom. Evaluate the Helmholtz free energy of the gas,  $F = -NT \log Z$ , where  $N$  is the total number of atoms in the gas, and find its equation of state using

$$P = -\left(\frac{\partial F}{\partial V}\right)_T$$

*Solution*

$$\begin{aligned} Z &= \frac{4\pi V}{(2\pi\hbar)^3} \int_0^\infty dp p^2 \exp\left(-\frac{p^2}{2mT}\right) \\ &= \frac{(2mT)^{3/2} V}{2\pi^2\hbar^3} \int_0^\infty dp p^2 \exp(-p^2) \\ &= -\frac{(2mT)^{3/2} V}{2\pi^2\hbar^3} \frac{\partial}{\partial \alpha} \int_0^\infty dp \exp(-\alpha p^2) \Big|_{\alpha=1} \\ &= -\frac{(2mT)^{3/2} V}{2\pi^2\hbar^3} \frac{\partial}{\partial \alpha} \left(\frac{1}{2}\sqrt{\frac{\pi}{\alpha}}\right) \Big|_{\alpha=1} \\ &= \frac{(2mT)^{3/2} V \sqrt{\pi}}{8\pi^2\hbar^3} = V \left(\frac{mT}{2\pi\hbar^2}\right)^{3/2} \end{aligned}$$

whereof

$$\begin{aligned} F &= -NT \log V \left(\frac{mT}{2\pi\hbar^2}\right)^{3/2} \\ P &= \frac{NT}{V} \end{aligned}$$

3. To find the Green's function of the Helmholtz equation in 3D, the following integral must be evaluated:

$$G_3(r) = -\frac{1}{(2\pi)^2 r} \int_{-\infty}^{\infty} dq \frac{q \sin qr}{q^2 - k^2}$$

Evaluate  $G_3(r)$  for the cases below (pick any two):

- (a) when both poles are moved slightly off the axis to the upper plane;
- (b) when the poles are moved slightly off the axis; the pole at  $k$  to the upper plane, and the pole  $-k$  to the lower plane;

- (c) when the poles are moved slightly off the axis; the pole at  $k$  to the lower plane, and the pole  $-k$  to the upper plane;  
 (d) when the integral is understood as the principal value.

*Solution*

(a)

$$G_3(r) = -\frac{1}{4\pi r} \left\{ \text{res} \left[ \frac{q [\exp(iqr)]}{q^2 - k^2}, k \right] + \text{res} \left[ \frac{q [\exp(iqr)]}{q^2 - k^2}, -k \right] \right\} = -\frac{\cos kr}{4\pi r}$$

(b)

$$\begin{aligned} G_3(r) &= -\frac{1}{4\pi r} \left\{ \text{res} \left[ \frac{q [\exp(iqr)]}{q^2 - k^2}, k \right] - \text{res} \left[ \frac{q [-\exp(-iqr)]}{q^2 - k^2}, -k \right] \right\} \\ &= -\frac{1}{4\pi r} \left[ \frac{\exp(ikr)}{2} + \frac{\exp(ikr)}{2} \right] = -\frac{\exp(ikr)}{4\pi r} \end{aligned}$$

(c)

$$\begin{aligned} G_3(r) &= -\frac{1}{4\pi r} \left\{ \text{res} \left[ \frac{q [\exp(iqr)]}{q^2 - k^2}, -k \right] - \text{res} \left[ \frac{q [-\exp(-iqr)]}{q^2 - k^2}, k \right] \right\} \\ &= -\frac{1}{4\pi r} \left[ \frac{\exp(-ikr)}{2} + \frac{\exp(-ikr)}{2} \right] = -\frac{\exp(-ikr)}{4\pi r} \end{aligned}$$

(d)

$$\begin{aligned} G_3(r) &= \frac{i}{2(2\pi)^2 r} \left[ P \int_C dq \frac{q \exp(iqr)}{q^2 - k^2} - P \int_C dq \frac{q \exp(-iqr)}{q^2 - k^2} \right] \\ &= \frac{i}{2(2\pi)^2 r} \left\{ \pi i \left( \text{res} \left[ \frac{q \exp(iqr)}{q^2 - k^2}, k \right] + \text{res} \left[ \frac{q \exp(iqr)}{q^2 - k^2}, -k \right] \right) \right\} \\ &\quad - \frac{i}{2(2\pi)^2 r} \left\{ -\pi i \left( \text{res} \left[ \frac{q \exp(-iqr)}{q^2 - k^2}, k \right] + \text{res} \left[ \frac{q \exp(-iqr)}{q^2 - k^2}, -k \right] \right) \right\} \\ &= -\frac{1}{8\pi r} \left( \frac{\exp(ikr)}{2} + \frac{\exp(-ikr)}{2} + \frac{\exp(-ikr)}{2} + \frac{\exp(ikr)}{2} \right) = -\frac{\cos kr}{4\pi r} \end{aligned}$$

4. Fresnel integrals, used in many fields of physics (e.g. optics), are defined as follows:

$$\begin{aligned} S(x) &= \int_0^x \sin\left(\frac{\pi t^2}{2}\right) dt \\ C(x) &= \int_0^x \cos\left(\frac{\pi t^2}{2}\right) dt \end{aligned}$$

Find their approximate expansions for small and large arguments. (Limit your expansions to  $O(x^3)$  for  $x \ll 1$  and  $O(x^{-1})$  for  $x \gg 1$ ). *Hint:* to evaluate  $S(\infty)$  and  $C(\infty)$ , evaluate the integral  $\oint \exp(-z^2/2) dz$  along the contour consisting of (i) the real axis from 0 to  $+\infty$ , (ii) one quarter of a large circle at  $|z| = \infty$ , and (iii) return to the origin along the line  $\arg(z) = \pi/4$ .

*Solution*

For  $x \ll 1$ ,

$$\begin{aligned} S(x) &= \int_0^x \sin\left(\frac{\pi t^2}{2}\right) dt \approx \int_0^x \frac{\pi t^2}{2} dt = \frac{\pi x^3}{6} \\ C(x) &= \int_0^x \cos\left(\frac{\pi t^2}{2}\right) dt \approx \int_0^x dt = x \end{aligned}$$

To evaluate  $S(\infty)$  and  $C(\infty)$ ,

$$\begin{aligned} 0 &= \oint \exp\left(-\frac{z^2}{2}\right) dz = \int_0^\infty \exp\left(-\frac{t^2}{2}\right) dt - \frac{1}{\sqrt{2}} \int_0^\infty \left[ \cos\left(\frac{t^2}{2}\right) + \sin\left(\frac{t^2}{2}\right) \right] dt \\ &\quad + \frac{i}{\sqrt{2}} \int_0^\infty \left[ \sin\left(\frac{t^2}{2}\right) - \cos\left(\frac{t^2}{2}\right) \right] dt \end{aligned}$$

and equating Re and Im parts to zero in the r.h.s.,

$$\int_0^\infty \sin\left(\frac{\pi t^2}{2}\right) dt = \int_0^\infty \cos\left(\frac{\pi t^2}{2}\right) dt = \frac{1}{2}$$

For  $x \gg 1$ ,

$$\begin{aligned} S(\infty) - S(x) &= \int_x^\infty \sin\left(\frac{\pi t^2}{2}\right) dt = -\frac{1}{\pi} \int_x^\infty \frac{1}{t} d \cos\left(\frac{\pi t^2}{2}\right) \approx \frac{\cos \pi x^2/2}{\pi x} \\ S(x) &\approx \frac{1}{2} - \frac{\cos \pi x^2/2}{\pi x} \end{aligned}$$

$$\begin{aligned} C(\infty) - C(x) &= \int_x^\infty \cos\left(\frac{\pi t^2}{2}\right) dt = \frac{1}{\pi} \int_x^\infty \frac{1}{t} d \sin\left(\frac{\pi t^2}{2}\right) \approx \frac{\sin \pi x^2/2}{\pi x} \\ C(x) &\approx \frac{1}{2} + \frac{\sin \pi x^2/2}{\pi x} \end{aligned}$$