Microwave Resonant Circuits and Structures. A.M. Ferendeci

In order to understand the properties of the microwave structures, well known series and parallel RLC circuits will be reviewed and their relation to microwave resonant circuits will be made.

Series RLC circuit



Consider first the series connected RLC circuit as shown in the figure. The input impedance can be written as

$$Z_{in} = R + j L + \frac{l}{j C}$$
(1)

This can be written as

$$Z_{in} = R + j L 1 - \frac{1}{{}^{2}LC}$$

$$Z_{in} = R + j L 1 - \frac{2}{{}^{0}}$$
(2)

Here $_{0}=1/\text{sqrt}(\text{LC})$ is the resonant frequency. Let

$$= _{0} +$$
(3)

where is the deviation of the frequency from resonant frequency $_{0}$. Substituting Eq.3 into Eq.2

$$Z_{in} = R + j(_{0} +) L 1 - \frac{2}{(_{0} +)^{2}}$$

$$Z_{in} = R + j(_{0} +) L 1 - \frac{2}{0}$$

$$Z_{in} = R + j(_{0} +) L 1 - \frac{1}{1 + 2 - \frac{1}{0} + \frac{2}{0}}$$

Neglecting second order terms and expanding 1/(1+x) as 1-x for x<<1 (x= / ____)

$$Z_{in} = R + j(_{o} +) L = 1 - 2 - \frac{1}{o}$$

$$Z_{in} = R + j(_{o} +) L = 2 - \frac{1}{o}$$

$$Z_{in} = R + j(2 - L)$$
(4)

The input impedance can be written in terms of power loss and energy stored

$$Z_{in} = \frac{P_{l} + 2j (W_{m} - W_{e})}{\frac{1}{2}II^{*}}$$
(5)
Where

$$P_1 = \frac{1}{2} R I I^*$$
 $W_m = \frac{1}{4} L I I^*$ $W_e = \frac{1}{4} \frac{I I^*}{C}$

The quality factor can be written as

$$Q = \frac{(\text{ total energy stored})}{\text{Power loss}} = \frac{\frac{1}{2} \frac{\text{I I}^*}{2\text{C}}}{\frac{1}{2} \text{ R I I}^*} = \frac{1}{2} \frac{1}{$$

Therefor as the series resistance decreases the quality factor increases. Eq.4 can now be written in terms of the Q of the circuits.

$$Z_{in} = R \quad 1 + j 2 - Q \tag{6}$$

The magnitude and phase of Eq.6 can now be plotted against frequency .



Note that at resonance when $=_{0}$, the input impedance magnitude is equal to R and the phase is 0° . Let $_{1} = _{0}$ - and $_{2} = _{0}$ + be the frequencies at which real and imaginary parts of Z_{in} are equal, i.e.,

$$\frac{2}{0} Q = \frac{2Q}{0} (0 - 1) = 1$$
$$1 = 0 \quad 1 - \frac{1}{2Q}$$

and

$$\frac{2}{0} Q = \frac{2Q}{0} (2 - 0) = 1$$
$$2 = 0 + \frac{1}{2Q}$$

the difference between 2^{-1} is defined as the bandwidth of the tuned circuit

$$_{2} - _{1} = _{0} 1 + \frac{1}{2Q} - _{0} 1 - \frac{1}{2Q} = _{0} \frac{1}{Q}$$

and the fractional bandwidth BW is defined as

$$BW = \frac{2^{-1}}{0} = \frac{1}{Q}$$

Note that the bandwidth is inversely proportional to the Q of the resonant structure. The higher the Q of the circuit, the narrower is the bandwidth of the circuit.

Parallel RLC circuit



The same procedure is used for the derivation of Z_{in} .

$$Z_{in} = \frac{1}{R} + j C - \frac{1}{j L}^{-1}$$

$$Z_{in} = \frac{1}{R} + j C 1 - \frac{1}{^{2}LC}^{-1}$$

$$Z_{in} = \frac{1}{R} + j C 1 - \frac{2}{^{0}}^{-1}$$

$$Z_{in} = \frac{1}{R} + j(_{0} + _{)} C 1 - \frac{2}{^{0}(_{(+)})^{2}}^{-1}$$

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$$Z_{in} = \frac{1}{R} + j(_{0} + _{)} C 1 - \frac{2}{^{0}(_{(+)})^{2}}^{-1}$$

$$Z_{in} = \frac{1}{R} + j(_{0}) C 2 - \frac{1}{^{0}}^{-1}$$

$$Z_{in} = \frac{1}{R} + j(_{0}) C 2 - \frac{1}{^{0}}^{-1}$$

$$Z_{in} = \frac{1}{R} + j C 2 - \frac{1}{^{0}}^{-1}$$

In terms of energies and power loss

$$P_{1} = \frac{1}{2} \frac{V^{2}}{R} \quad W_{e} = \frac{1}{4}C \quad V^{2} \quad W_{m} = \frac{1}{4}L \quad I^{2} = \frac{1}{4}L \frac{V}{jL}^{2} = \frac{1}{4}\frac{V^{2}}{_{o}L}$$

$$Q = \frac{\text{(Total energy stored)}}{\text{Total power loss}} = \frac{\frac{\circ}{2}C \quad V^{2}}{\frac{1}{2}\frac{V^{2}}{R}} = \frac{\circ CR}{_{o}CR} = \frac{R}{_{o}L}$$

In the case of parallel circuit, the Q increases with increasing R. $Z_{\mbox{\tiny in}}$ can be written in terms of Q as

$$Z_{in} = \frac{1}{(G + j2 - C)} = \frac{R}{1 + j2 - RC} = \frac{R}{1 + jQ - C}$$

Magnitude and phase of Z_{in} as a function of for the parallel RLC is plotted in the figure below.



Loaded Q

The circuits given above are isolated circuit elements. In actual circuitry, the resonant structure is connected to the outside world or signal is coupled from an external system. The resonant circuit is usually connected to a transmission line and thus sees at its terminals equivalent resistance R_e .

For the Series circuit, the external resistance is in series with the R of the isolated system. The total resistance is then

$$R_{t} = R + R_{e}$$



The equivalent Q of the system including the external resistance is now defined as the loaded Q, or Q_L .

$$Q_{L} = \frac{1}{{_{o}CR_{t}}} = \frac{1}{{_{o}C(R_{e}+R)}} = \frac{1}{{_{o}CR_{e}} + {_{o}CR}} = \frac{1}{\frac{1}{\frac{1}{Q_{e}} + \frac{1}{Q_{o}}}}$$
$$\frac{1}{Q_{L}} = \frac{1}{Q_{e}} + \frac{1}{Q_{o}}$$

Here $Q_o = 1/CR$ is the unloaded Q and the $Q_e = 1/CR_e$ is the external Q of the cavity. The measurement provides the loaded Q_L which is always smaller than the unloaded Q_o of the resonant structure. For the parallel circuitry, the external resistance is in parallel with the isolated

resistance.



The total resistance is $R_t = [RR_e/(R+R_e)]$. The loaded Q for this case is

$$Q_{L} = {}_{o}C R_{t} = {}_{o}C \frac{R_{e}R}{R_{e}+R} = {}_{o}C \frac{1}{\frac{1}{R_{e}} + \frac{1}{R}} = \frac{1}{\frac{1}{_{o}CR_{e}} + \frac{1}{_{o}CR}}$$
$$Q_{L} = \frac{1}{\frac{1}{Q_{e}} + \frac{1}{Q_{o}}}$$
$$\frac{1}{Q_{L}} = \frac{1}{Q_{e}} + \frac{1}{Q_{o}}$$

This is the same result as the Series case.