



ELSEVIER

Available online at www.sciencedirect.com

SCIENCE @ DIRECT®

Physica C xxx (2003) xxx–xxx

PHYSICA C

www.elsevier.com/locate/physc

Suppression of *ab*-plane crack formation in single domain $\text{YBa}_2\text{Cu}_3\text{O}_x$ by uniaxial *c*-axis pressure

D. Shi ^{a,*}, D. Isfort ^{b,c}, X. Chaud ^c, P. Odier ^b, D. Mast ^d, R. Tournier ^{b,c}

^a Department of Chemical and Materials Engineering, University of Cincinnati, Cincinnati, OH 45221-0012, USA

^b Laboratoire de Cristallographie, CNRS, BP 166, F-38042 Grenoble Cedex 9, France

^c Consortium de Recherches pour l'Emergence de Technologies Avancées, CNRS, BP 166, F-38042 Grenoble Cedex 9, France

^d Department of Physics, University of Cincinnati, Cincinnati, OH 45221-0012, USA

Received 19 May 2003; received in revised form 10 July 2003; accepted 25 August 2003

Abstract

A uniaxial *c*-axis pressure has been applied to $\text{YBa}_2\text{Cu}_3\text{O}_x$ (YBCO) single domain during a prolonged oxygenation experiment up to 72 h. It has been found that, due to this *c*-axis pressure, the local tensile stress responsible for *ab*-plane crack formation can be compensated effectively. The optical microscopy experimental results indicate that there is a sharp contrast in microstructure between the samples with and without *c*-axis pressure. The *ab*-plane cracks in the sample with *c*-axis pressure appear to be suppressed. X-ray diffraction results show that, for a 72 h-oxygen anneal, both samples with and without pressure are well oxygenated with equal orthorhombicity on the polished surfaces where the microstructures are studied. Also discussed is the mechanism of *ab*-plane crack suppression by the *c*-axis pressure.

© 2003 Elsevier B.V. All rights reserved.

1. Introduction

In $\text{YBa}_2\text{Cu}_3\text{O}_x$ (YBCO) single domain, *ab*-plane cracks inevitably form due to large stresses induced during the tetragonal (T) to orthorhombic (O) phase transformation [1–3]. Therefore, crack formation is inherently a structure-related behavior governed by oxygen diffusion. With these cracks in the matrix of the crystal, the integrity of the domain structure is severely altered making its usefulness limited. For instance, the mechanical strength of the YBCO single domain is signifi-

cantly reduced and the materials easily cleave along the *ab*-plane cracks. In magnetic levitation, these cracks interfere with the induced currents; thus lower the trapped field and levitation force. In fundamental studies, these imperfections make it difficult to control the physical parameters that are needed for modeling.

One of the novel applications is in the RF wireless telecommunications. The recent development of extremely low loss RF components using high temperature superconductors (HTS) such as YBCO has succeeded in achieving the extremely low surface resistance at the appropriate cellular frequencies. The current technologies in fabrication of HTSs for RF applications have been primarily the thin film approaches that can produce

* Corresponding author. Tel.: +1-513-556-3100; fax: +1-513-556-1004.

E-mail address: donglu.shi@uc.edu (D. Shi).

47 well-textured $\text{YBa}_2\text{Cu}_3\text{O}_x$ materials with extremely
48 low surface resistance [4,5]. The recent develop-
49 ment in seeded melt growth (SMG) of YBCO has
50 shown great promise in the fabrication of RF
51 components with superb RF properties. Not only
52 has the surface resistance of the SMG reached a
53 value comparable to those of thin films, great ad-
54 vantages in utilizing this HTS material also lie in
55 its suitability for mass production, easy control of
56 device geometry, and low cost.

57 One of the concerns for single domain YBCO is
58 that it inherently contains these *ab*-plane micro-
59 cracks. As RF waves travel through the surfaces,
60 great loss can take place at these cracked regions.
61 Note that the microwaves only penetrate about
62 $1\mu\text{m}$ depth on the surface, therefore these surface
63 defects become more critical in reducing the RF
64 losses and device design. There have been, how-
65 ever, a few systematic studies so far to investigate
66 the effective methods that can prevent these sur-
67 face cracks.

68 Suasmoro was the first to use in situ ultrasonic
69 wave attenuation to probe the microstructure of
70 YBCO ceramic polycrystals [6]. He concluded,
71 based on the ultrasonic data, that the microcracks
72 formed in the tetragonal phase region where oxi-
73 dation was taking place. Diko had made an ex-
74 tensive analysis on the behavior of microcracks in
75 textured YBCO [7]. He discussed on the origin of
76 these cracks: i.e., whether these are growth defects
77 or oxygen annealing defects. Growth defects were
78 found parallel to the *ab*-planes due to a growth
79 gap generated by inclusions of 211 particles dur-
80 ing crystallization. Diko attributed the micro-
81 cracking phenomena to a thermal expansion stress
82 between oxidized and non-oxidized zones. The
83 oxygenation was thought to be a combination of
84 oxygen volume diffusion, microcracking, and ox-
85 ygen penetration along the cracks.

86 In a recent study by Shi et al. on crack forma-
87 tion and propagation in YBCO single domain,
88 they concluded that, during oxygenation, the
89 phase transition has created a large volume stress
90 that is responsible for rapid crack propagation in
91 the YBCO crystal [8]. Due to severe contraction in
92 the *c*-direction, the crystal under goes a compres-
93 sive deformation. However the structurally iso-
94 tropic 211 particles prevent the compressive

95 deformation in some local areas especially near the
96 211 particles, therefore diverting the compressive
97 stress to tensile stress. As the tensile stress reaches
98 the critical value, the cracks begin to propagate
99 inwardly in the crystal. The T-to-O phase bound-
100 ary provides the driving force for the propagation
101 of the cracks. Since the T-to-O transition is oxygen
102 diffusion controlled, the kinetics of crack propa-
103 gation is dictated by the rate of oxygen diffusion.
104 After the oxygen content is homogeneously es-
105 tablished in the sample, there is no further driving
106 force for the propagation of cracks.

107 Based on these studies, in this paper, we report
108 on a new method by which the *ab*-plane crack
109 formation was suppressed by applying a uniaxial
110 *c*-axis pressure to the single domain directly. Both
111 optical microscopy and X-ray diffraction results of
112 the annealed samples are presented. The mecha-
113 nism on suppression of crack formation is dis-
114 cussed.

2. Experimental

115
116 The seeded melt growth (SMG) method was
117 originally developed for the purpose of producing
118 large domains for various engineering applications
119 [9–12]. The single domain samples were disc
120 shaped pellets (20 mm in diameter and 15 mm in
121 height). We used precursors of 70 wt%
122 $\text{YBa}_2\text{Cu}_3\text{O}_x$, 30 wt% Y_2BaCuO_5 , and 0.15 wt% Pt.
123 These powders were thoroughly mixed and
124 uniaxially pressed at 100 MPa into a disc shaped
125 green pellet of 35 g. The green pellet was sintered
126 in air at 930 °C for 24 h. A $\text{SmBa}_2\text{Cu}_3\text{O}_x$ single
127 domain was used as seed having a dimension of
128 $2 \times 2 \times 1.5 \text{ mm}^3$. The seed was put on the top of the
129 green pellet before the growth process. The green
130 pellet was placed on an alumina plate with an in-
131 termediate layer of Y_2O_3 powder to avoid liquid
132 spreading on the interior surfaces of the furnace. A
133 sintered thin plate consisting of a mixture of
134 $\text{YbBa}_2\text{Cu}_3\text{O}_x$ and $\text{YBa}_2\text{Cu}_3\text{O}_x$ was also used be-
135 tween the green pellet and the Y_2O_3 layer. This
136 substrate was found necessary to avoid parasitic
137 grain growth from the bottom. After the comple-
138 tion of the domain growth, the sample was cooled
139 in nitrogen gas to prevent the phase transforma-

140 tion from the tetragonal to orthorhombic struc-
 141 ture. In this way, the sample is free of *ab*-plane
 142 crack formation. The details about the YBCO
 143 single domain growth can be found in Refs. [13–
 144 16].

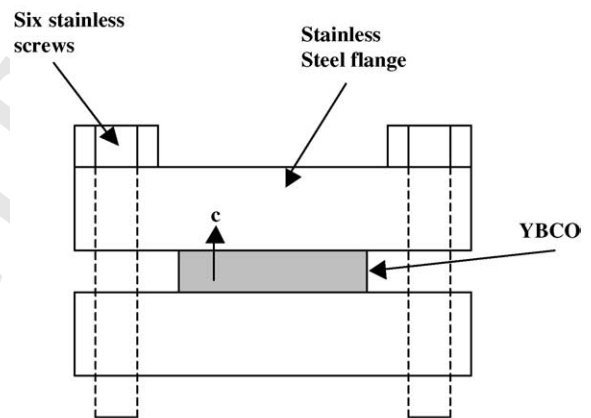
145 In this experiment, a 2.5-mm-thick YBCO disc
 146 of 20 mm diameter was sliced along the *ab*-plane
 147 from the crack-free single domain pellet for the
 148 pressure experiment. The virgin sample had the
 149 tetragonal phase as it was cooled in nitrogen gas.
 150 The disc sample was cut, parallel to the *c*-axis, into
 151 two pieces with the facing surfaces polished down
 152 to 1 μm for optical microscopy studies. The mir-
 153 ror-image surfaces were used to ensure that the
 154 studied regions were adjacent to each other. The
 155 pressure vise was a commercial flange made of
 156 stainless steel as shown in Fig. 1. The flange di-
 157 ameter was 34 mm with a thickness of 7.2 mm. As
 158 shown in Fig. 1b, the disc sample was sandwiched
 159 between two flanges tightened with screws. Before
 160 tightening the screws, the flanges were placed in a
 161 hydraulic press for achieving a uniform force
 162 along *c*-axis. As the hydraulic pressure reached 15
 163 MPa, six 4.2 mm screws were hand-tightened. To
 164 ensure a uniform pressure, the top and bottom
 165 surfaces of the disc sample were polished by a
 166 special device making these surfaces parallel to
 167 each other. Note that a device that has resulted in
 168 a more uniform force on Bi2223 has been reported
 169 in a previous work [17]. To study the effect of
 170 pressure on the crack behavior, we are in the
 171 process of designing a better device and monitor
 172 the pressure at the annealing temperatures.

173 Both samples (each half) with and without
 174 pressure were annealed together in flowing oxygen
 175 at 400 $^{\circ}\text{C}$ for 18, 48 and 72 h. At each annealing
 176 interval, the samples were cooled to room tem-
 177 perature and used for optical microscopy and X-
 178 ray diffraction (XRD). After these studies, the
 179 samples were put back into the furnace for a
 180 prolonged oxygenation under the same condition
 181 up to 72 h.

182 A thin layer from the polished surfaces of the
 183 oxygenated samples (with and without pressure)
 184 was powdered and analyzed by X-ray diffraction
 185 (XRD) using a Siemens diffractometer (D5000)
 186 operating in transmission mode with Cu $K\alpha$ radi-
 187 ation.



(a)



(b)

Fig. 1. (a) A photograph showing the stainless steel pressure vise and (b) schematic diagram showing the application of uniaxial *c*-axis pressure.

3. Results

188

189 Fig. 2a shows the microstructure of the as-
 190 grown YBCO single domain. Due to flowing ni-
 191 trogen during cooling, the YBCO phase remains
 192 un-oxygenated with a tetragonal structure. No
 193 cracks are observed in this sample. However, after
 194 the sample was treated in flowing oxygen at 400 $^{\circ}\text{C}$
 195 for 18 h, microcracks formed as shown in Fig. 2b,
 196 which have been typically seen in many previous

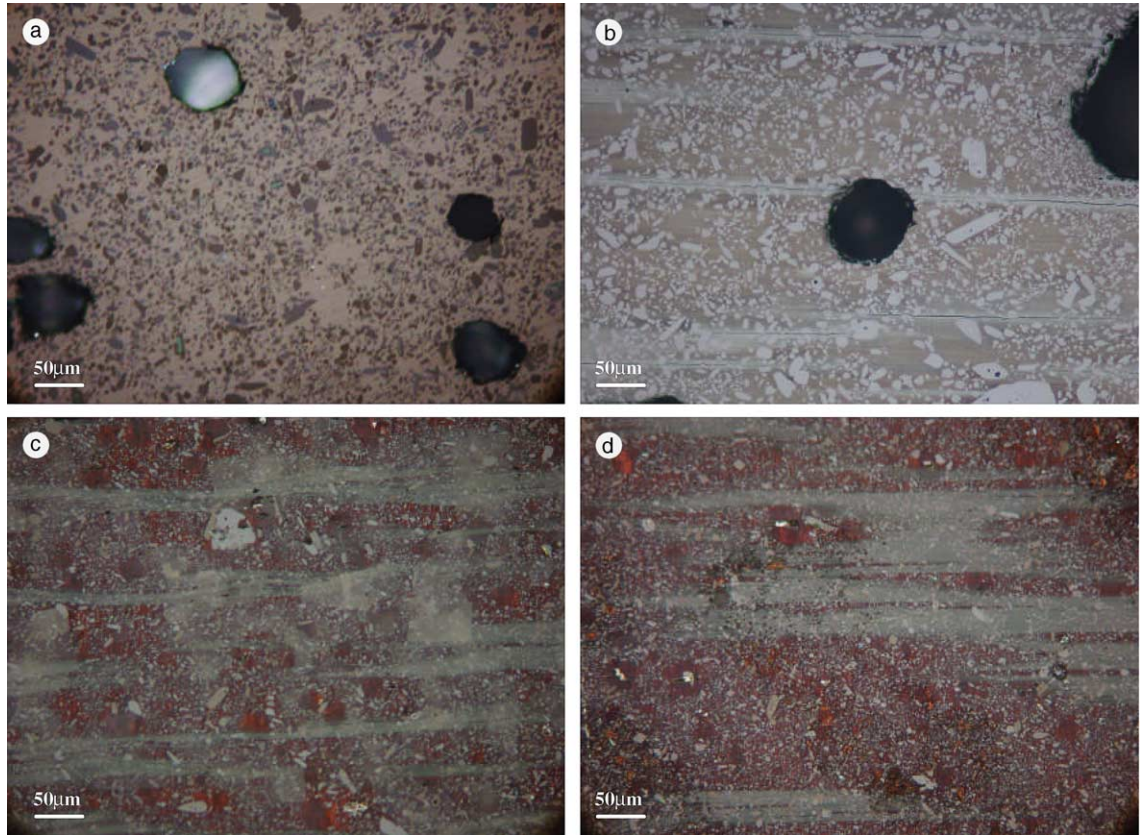


Fig. 2. (a) Optical micrograph showing the microstructure of the as-grown single domain, (b) optical micrograph showing the microstructure of the sample annealed at 400 °C for 18 h without pressure, (c) optical micrograph showing the microstructure of the sample annealed at 400 °C for 18 h with the *c*-axis pressure, and (d) the same as in (c) but another area.

197 studies. Note that there are regions which have
 198 different contrast with the matrix in which the
 199 cracks are visible. These regions are assumed to be
 200 more oxygenated than the matrix. Upon crack
 201 formation, oxygen diffusion becomes more rapid
 202 along the crack. As the oxygen diffuses into the
 203 sample along the *ab*-planes, these microcracks also
 204 propagate into the interior of the domain structure
 205 at a rate that is comparable to that of oxygen
 206 diffusion. However, in this study we only focused
 207 on the *polished surfaces* of the annealed samples
 208 with and without pressure. The crack formation
 209 inside the matrix of the sample is being investi-
 210 gated in a more detailed study. Fig. 2c shows the
 211 counterpart of the sample that was annealed under
 212 the *c*-axis pressure. Regions with light contrast,
 213 similar to those shown in Fig. 2b, are observed. In

214 these regions, no visible cracks were observed with
 215 optical microscopy. Fig. 2d shows another area of
 216 the sample. Again it indicates a quite homoge-
 217 neous microstructure throughout the entire sam-
 218 ple.

219 Fig. 3 shows the microstructures of the samples
 220 annealed for 48 h with and without *c*-axis pressure.
 221 Quite similar microstructural behaviors shown in
 222 Fig. 2 have been observed in these pictures (e.g.
 223 with *c*-axis pressure, no microcracks are observed).
 224 We found the identical microstructural differences
 225 between the samples with and without *c*-axis
 226 pressure as shown in Fig. 4. Up to an annealing
 227 time of 72 h, the sample with pressure remains
 228 uncracked as can be seen in this figure. As these
 229 *ab*-plane cracks were initiated by oxygen diffusion,
 230 for samples with *c*-axis pressure applied we needed

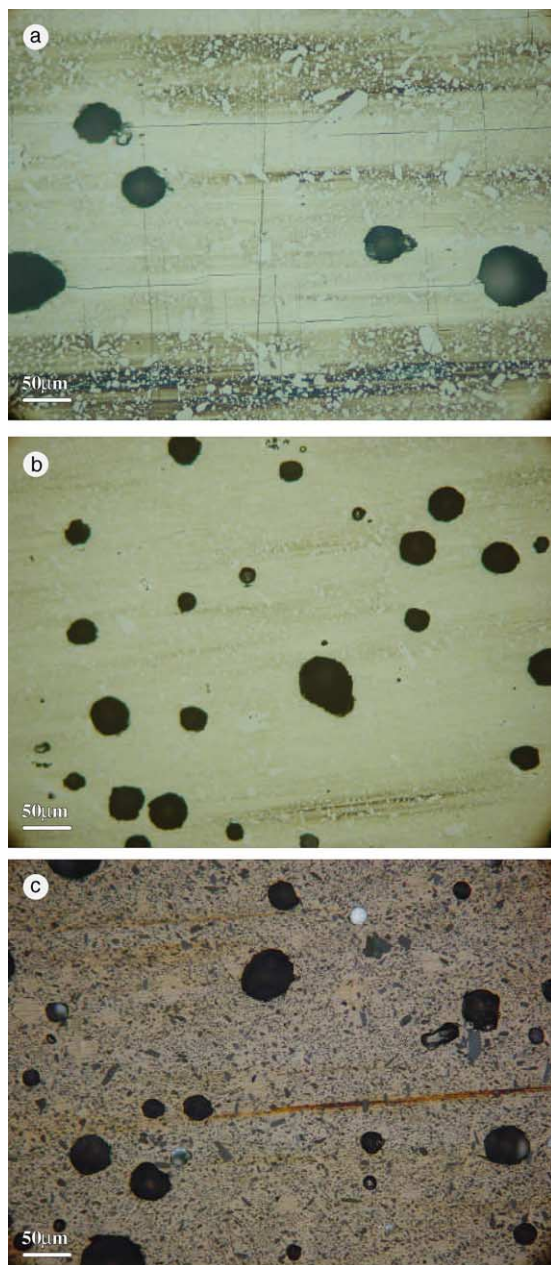


Fig. 3. (a) Optical micrograph showing the microstructure of the sample annealed at 400 °C for 48 h without pressure, (b) optical micrograph showing the microstructure of the sample annealed at 400 °C for 48 h with the *c*-axis pressure, and (c) the same as in (b) but another area.

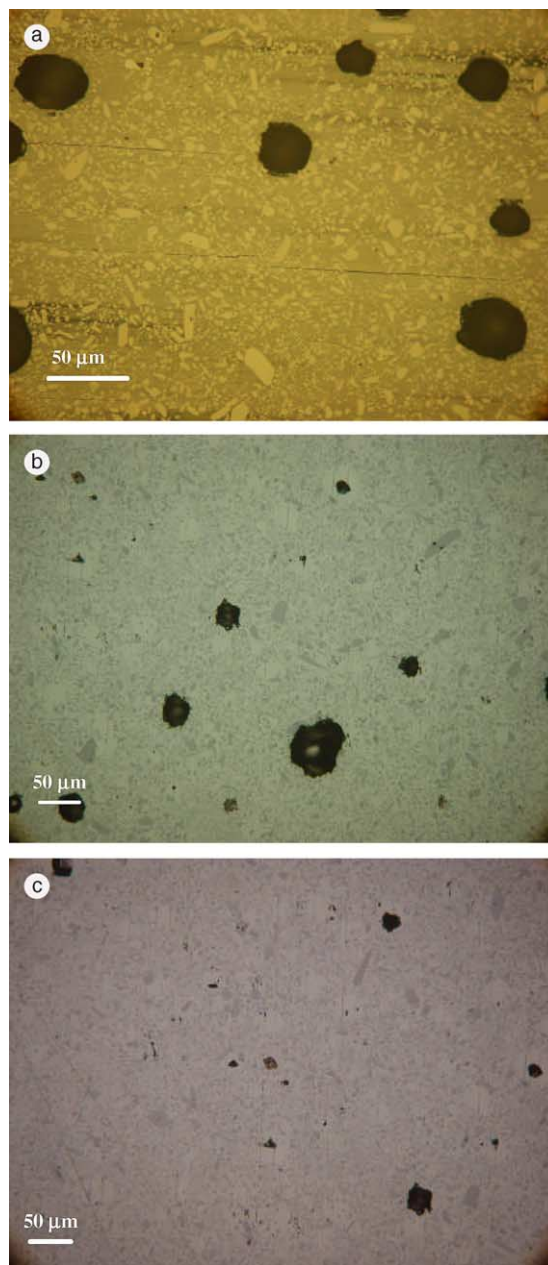


Fig. 4. (a) Optical micrograph showing the microstructure of the sample annealed at 400 °C for 72 h without pressure; (b) optical micrograph showing the microstructure of the sample annealed at 400 °C for 72 h with the *c*-axis pressure, and (c) the same as in (b) but another area.

231 to examine the oxygen content in these samples
232 with and without pressure.

To determine the oxygen content and the orthorhombicity of the annealed samples, XRD was

233
234

235 carried out on both polished surfaces for the 72 h-
 236 annealed samples with and without pressure. Small
 237 portions of material were taken from these sample
 238 surfaces (i.e. the sample with pressure and the
 239 sample without pressure) for the XRD powder
 240 diffraction. The spectrum was recorded in the angular
 241 range 10–100° in 2θ and the lines of
 242 $\text{YBa}_2\text{Cu}_3\text{O}_x$ phase were indexed according to the
 243 Pmmm space group. Only the lines that did not
 244 overlap with those of the Y_2BaCuO_5 phase (2 1 1)
 245 were taken into account for calculation of the unit
 246 cell parameters by a least square refinement
 247 method (Cellref program). We then used 25 lines,
 248 which were sufficient to provide accurate indexing.
 249 In a second attempt we mounted the 72 h-annealed
 250 samples directly on a four-circle goniometer. We
 251 registered the first $h00$ and $0k0$ reflections and
 252 used the Nelson–Riley procedure to extrapolate a
 253 and b values [18]. In this procedure, the calculated
 254 parameters for each set of data were plotted versus
 255 $\cos^2\theta/\sin\theta$ (Fig. 5). The extrapolation to zero
 256 gives a very good approximation of the cell param-
 257 eter. Both methods agree to point out that
 258 these samples have the identical lattice parameters
 259 (see Fig. 5a and b), indicating that they are fully
 260 oxygenated and have an orthorhombic structure at
 261 the polished surfaces.

262 **4. Discussion**

263 Applications of pressure on oxide supercon-
 264 ductors have been reported in the past for
 265 achieving various unique properties. Zhu et al.
 266 have reported a compressive anneal processing of
 267 Bi2223 superconducting tapes for enhancement of
 268 texture [17]. During oxygenation, tetragonal to
 269 orthorhombic phase transition takes place that
 270 generates volume stresses due to dimensional
 271 changes of the unit cell. To accommodate the
 272 volume stresses, twinning occurs as have been well
 273 observed in previous studies [1–4]. To study the
 274 effects of twins on various physical properties such
 275 as flux pinning and superconducting fluctuations,
 276 moderate ab -plane uniaxial pressure was applied
 277 to remove the twin boundaries [19–21]. However,
 278 we believe that, it is the first time that a c -axis

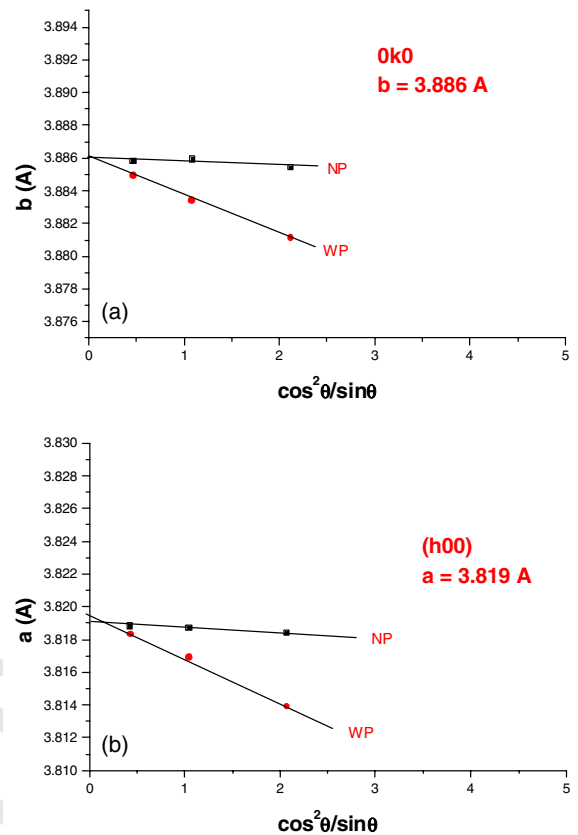


Fig. 5. (a) Lattice parameter b versus $\cos^2\theta/\sin\theta$ for both samples with and without pressure annealed at 400 °C for 72 h and (b) Lattice parameter a versus $\cos^2\theta/\sin\theta$ for both samples with and without pressure annealed at 400 °C for 72 h.

pressure has been applied to suppress ab -plane crack initiation and propagation.

The ab -plane crack formation has extensively been studied previously [22,23]. During oxygen diffusion from the surface of the YBCO sample, an oxygen gradient is created. As a result of c -axis contraction, a tensile stress is created in the outer regions of the sample. This strain profile in the oxygenated layer will cause cracking parallel to ab -plane with cracks apparent at the sample surface [22]. The 2 1 1 particles in the YBCO single domain can also contribute to the ab -plane cracks [23]. This local tensile stress will initiate cracks along ab -planes and the level of this stress depends on the size of 2 1 1 particles, presumably, the larger the particle, the higher the stress. From our mi-

279
 280
 281
 282
 283
 284
 285
 286
 287
 288
 289
 290
 291
 292
 293
 294

295 crostructural observation the particles size range
 296 from submicron up to 20 μ with an average size of
 297 2 μ m and a few large 2 1 1 particles resulted from
 298 peritectic reaction, therefore can generate stresses
 299 in a wide range. As the YBCO single domain
 300 contains a high concentration of 2 1 1 particles (30
 301 wt% in initial concentration), they will hinder the
 302 *c*-axis change and create a local tensile stress near
 303 themselves. The *c*-axis variation with annealing
 304 time was previously published by Shi et al. [24].
 305 Such *ab*-plane cracks near the 2 1 1 particles have
 306 been commonly observed [22–24]. Therefore to
 307 compensate this local tensile stress that is respon-
 308 sible for *ab*-plane cracks, a uniaxial *c*-axis pressure
 309 can be applied by using the simple vise shown in
 310 Fig. 1. As the vise applies the pressure along the *c*-
 311 axis, the single domain will deform elastically al-
 312 though moderately. The *c*-axis pressure would
 313 confine the crystal deformation along the *c*-axis
 314 during oxygenation. For a crack to initiate, local
 315 tensile stress is required. The *c*-axis pressure can
 316 provide an effective counter pressure that com-
 317 pensates this local tensile stress. In this way, the
 318 crack formation is suppressed. However, in this
 319 initial experiment, it is difficult to quantitatively
 320 determine the local crack tip stress and the counter
 321 force from the uniaxial pressure applied. More
 322 detailed experiments addressing these critical
 323 stresses are underway.

324 The *ab*-plane crack nucleation and propagation
 325 have been previously studied by Shi et al. [24]. In
 326 their experiments, Young's modulus experiments
 327 were carried out to study the kinetics of *ab*-plane
 328 crack propagation in single domain $\text{YBa}_2\text{Cu}_3\text{O}_x$
 329 (YBCO) during a prolonged oxygen heat treat-
 330 ment at 400 °C up to 188 h. It was found that the
 331 modulus value experiences a rapid fall between the
 332 annealing time of 48 and 96 h, indicating the ki-
 333 netics of crack propagation in the sample matrix.
 334 For 72 h annealed the sample, based on the results
 335 of this previous work, the *ab*-plane cracks have
 336 propagated through the entire cross section along
 337 the *ab*-plane. With the confinement of the *c*-axis
 338 pressure, not only can the crack initiation be sig-
 339 nificantly repressed, but also the kinetics slowed
 340 down. However, in this experiment, our experi-
 341 ments are only focused at the polished surfaces of
 342 annealed samples with and without *c*-axis pres-

sure. The effects of *c*-axis pressure on crack
 propagation in the matrix of the single domain are
 currently being carried out.

5. Summary

In order to suppress *ab*-plane crack formation,
 uniaxial *c*-axis pressure is applied to the YBCO
 single domain. In a prolonged oxygen anneal, this
c-axis pressure has effectively suppressed the for-
 mation of the *ab*-plane cracks on the *polished*
surfaces. These results indicate that by applying
 pressure along the *c*-axis, local tensile stresses re-
 sponsible for initiating *ab*-plane cracks can be
 compensated effectively; thereby providing a un-
 ique method for eliminating cracks in the YBCO
 single domain.

Acknowledgement

One the of authors (D. Shi) is grateful to the
 financial and technical support from Laboratoire
 de Cristallographie and Consortium de Recher-
 ches pour l'Emergence de Technologies Avancées,
 CNRS, Grenoble, France.

References

- [1] D. Shi, D.W. Capone II, K.C. Goretta, K. Zhang, G.T. Goudey, *J. Appl. Phys.* 63 (1988) 5411.
- [2] D. Shi, K. Zhang, D.W. Capone II, *J. Appl. Phys.* 64 (1995) 1988.
- [3] D. Shi, *Phys. Rev. B* 39 (1989) 4299.
- [4] J.D. Jorgensen, *Jpn. J. Appl. Phys.* 26 (1987) 2017.
- [5] T. Siegrist, S. Sunshine, D.W. Murphy, R.J. Cava, S.M. Zahurak, *Phys. Rev. B* 35 (1987) 7137.
- [6] Suasmoro, Ph.D. dissertation, 1992.
- [7] P. Diko, N. Pelerin, P. Odier, *Physica C* 247 (1995) 169.
- [8] D. Shi, P. Odier, A. Sulpice, D. Isfort, X. Chaud, R. Tournier, P. He, R. Singh, *Physica C* 384 (2003) 149.
- [9] D. Shi, W. Zhong, U. Welp, S. Sengupta, V. Todt, G.W. Crabtree, S. Dorris, U. Balachandran, *IEEE Trans. Magn.* 5 (1994) 1627.
- [10] M. Murakami, M. Morita, N. Koyama, *Jpn. J. Appl. Phys.* 28 (1989) L1125.
- [11] V.R. Todt, S. Sengupta, D. Shi, J. Hull, P.R. Sahm, P.J. McGinn, R. Peoppel, *J. Electron. Mater.* 23 (1994) 1127.
- [12] D. Shi et al., *Physica C* 246 (1995) 253.

- 385 [13] P. Gautier-Picard, X. Chaud, E. Beaugnon, A. Erraud, R. 396
386 Tournier, *Mater. Sci. Eng. B* 53 (1998) 66. 397
387 [14] P. Gautier-Picard, E. Beaugnon, X. Chaud, A. Sulpice, R. 398
388 Tournier, *Physica C* 308 (1998) 161. 399
389 [15] X. Chaud, D. Isfort, E. Beaugnon, R. Tournier, *Physica C* 400
390 2413 (2000) 341. 401
391 [16] K. Conder, *Mater. Sci. Eng.* 32 (2001) 41. 402
392 [17] Zhu et al., *Supercond. Sci. Technol.* 12 (1999) 640. 403
393 [18] H.P. Klug, L.E. Alexander, *X-ray Diffraction Procedures* 404
394 for Polycrystalline and Amorphous Materials, second ed., 405
395 John Wiley and Sons, 1974. 406
[19] R. Liang, D.A. Bonn, W.N. Hardy, *Phys. Rev. Lett.* 76 396
(1996) 835. 397
[20] M. Pissas, E. Moraitakis, G. Kallias, *Phys. Rev. B.* 62 398
(2000) 1446. 399
[21] U. Welp, J.A. Fendrich, W.K. Kwok, G.W. Crabtree, 400
B.W. Veal, *Phys. Rev. Lett.* 76 (1996) 4809. 401
[22] P. Diko, G. Krabbes, *Supercond. Sci. Technol.* 16 (2003) 402
90. 403
[23] D. Isfort, X. Chaud, R. Tournier, G. Kapelski, *Physica C* 404
390 (2003) 341. 405
[24] D. Shi, P. Odier, A. Sulpice, D. Isfort, X. Chaud, R. 406
Tournier, P. He, R. Singh, *Physica C* 384 (2003) 149. 407