Nonlinear oscillators at resonance with periodic forcing

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Abstract

In this note we unify the results of A.C. Lazer and P.O. Frederickson [3], A.C. Lazer [6], A.C. Lazer and D.E. Leach [7], J.M. Alonso and R. Ortega [1], and P. Korman and Y. Li [4] on periodic oscillations and unbounded solutions of nonlinear equations with linear part at resonance and periodic forcing. We give conditions for the existence and non-existence of periodic solutions, and obtain a rather detailed description of the dynamics for nonlinear oscillations at resonance, in case periodic solutions do not exist.

Key words: Resonance, periodic oscillations, unbounded solutions.

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1 Introduction

We are interested in the existence of 2π periodic solutions to the problem (here x = x(t))

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(1.1)
$$x'' + f(x)x' + g(x) + n^2 x = e(t)$$

The linear part

$$x'' + n^2 x = e(t)$$

is at resonance, with the null space spanned by $\cos nt$ and $\sin nt$. Define $F(x) = \int_0^x f(z) dz$. We assume throughout this paper that $e(t) \in C(R)$ satisfies $e(t + 2\pi) = e(t)$ for all t, $f(x), g(x) \in C(R)$, $n \ge 1$ is an integer; moreover, we assume that the finite limits at infinity $F(\infty)$, $F(-\infty)$, $g(\infty)$, $g(-\infty)$ exist, and

(1.2)
$$F(-\infty) < F(x) < F(\infty), \quad \text{for all } x \in R,$$

(1.3)
$$g(-\infty) < g(x) < g(\infty)$$
, for all $x \in R$.

Define

$$A_n = \int_0^{2\pi} e(t) \cos nt \, dt, \ B_n = \int_0^{2\pi} e(t) \sin nt \, dt$$

In the case when f = 0, the equation

(1.4)
$$x'' + g(x) + n^2 x = e(t)$$

was considered in the paper of A.C. Lazer and D.E. Leach [7] who proved the following classical theorem.

Theorem 1.1 ([7]) The condition

$$\sqrt{A_n^2 + B_n^2} < 2\left(g(\infty) - g(-\infty)\right)$$

is necessary and sufficient for the existence of 2π periodic solutions of (1.4).

In case g = 0, for the corresponding equation

(1.5)
$$x'' + f(x)x' + n^2x = e(t)$$

one has the following result.

Theorem 1.2 The condition

$$\sqrt{A_n^2 + B_n^2} < 2n \left(F(\infty) - F(-\infty) \right)$$

is necessary and sufficient for the existence of 2π periodic solutions of (1.5).

This theorem was proved by A.C. Lazer [6] for n = 1 (with an earlier result by P.O. Frederickson and A.C. Lazer [3]), and for all $n \ge 1$ by P. Korman and Y. Li [4], who used a small modification of A.C. Lazer's proof.

Question: is it possible to combine these theorems for the equation (1.1)? It turns out that the necessary conditions can be combined, while sufficient conditions cannot be combined.

Proposition 1 The condition

(1.6)
$$\sqrt{A_n^2 + B_n^2} < 2n \left(F(\infty) - F(-\infty)\right) + 2 \left(g(\infty) - g(-\infty)\right)$$

is necessary for the existence of 2π periodic solution of (1.1).

Proposition 2 The condition (1.6) is not sufficient for the existence of 2π periodic solution of (1.1).

In case (1.4) has no 2π periodic solutions, all solutions of (1.4) are unbounded as $t \to \pm \infty$, as follows by the second Massera's theorem, as was observed first by G. Seifert [8]. Later J.M. Alonso and R. Ortega [1] gave an elementary approach to this result (with a more refined statement, asserting that solutions tend to infinity in C^1 norm). We observe next that the approach of [1] works for the equation (1.1) as well.

Proposition 3 Assume in addition to the assumptions above that f(x) is uniformly bounded from below (the assumption (2.10) below). Then in case

(1.7)
$$\sqrt{A_n^2 + B_n^2} \ge 2n \left(F(\infty) - F(-\infty)\right) + 2 \left(g(\infty) - g(-\infty)\right),$$

all solutions of (1.1) satisfy $\lim_{t\to\pm\infty} \left(x^2(t) + {x'}^2(t)\right) = \infty$.

Clearly, there are no 2π periodic solutions in this case, in view of Proposition 1. Proposition 3 shows that the absence of 2π periodic solutions turns out to be more decisive in determining the overall dynamics of (1.1) than the existence of 2π periodic solutions.

2 The proofs

The following elementary lemmas are easy to prove.

Lemma 2.1 Consider a function $\cos(nt - \varphi)$, with an integer n and any real φ . Denote $P_c = \{t \in (0, 2\pi) \mid \cos(nt - \varphi) > 0\}$ and $N_c = \{t \in (0, 2\pi) \mid \cos(nt - \varphi) < 0\}$. Then

$$\int_{P_c} \cos(nt - \varphi) \, dt = 2, \quad \int_{N_c} \cos(nt - \varphi) \, dt = -2.$$

Lemma 2.2 Consider a function $\sin(nt - \varphi)$, with an integer n and any real φ . Denote $P_s = \{t \in (0, 2\pi) \mid \sin(nt - \varphi) > 0\}$ and $N_s = \{t \in (0, 2\pi) \mid \sin(nt - \varphi) < 0\}$. Then

$$\int_{P_s} \sin(nt - \varphi) \, dt = 2, \quad \int_{N_s} \sin(nt - \varphi) \, dt = -2 \, .$$

Proof of Proposition 1. Given arbitrary numbers *a* and *b*, one can find a $\delta \in [0, 2\pi)$, so that

$$a\cos nt + b\sin nt = \sqrt{a^2 + b^2}\cos(nt - \delta)$$
.

(with $\cos \delta = \frac{a}{\sqrt{a^2+b^2}}$, $\sin \delta = \frac{b}{\sqrt{a^2+b^2}}$.) It follows that

(2.1)
$$\frac{A_n}{\sqrt{A_n^2 + B_n^2}} \cos nt + \frac{B_n}{\sqrt{A_n^2 + B_n^2}} \sin nt = \cos(nt - \delta),$$

for some $\delta \in [0, 2\pi)$. Multiply (1.1) by $\frac{A_n}{\sqrt{A_n^2 + B_n^2}} \cos nt$ and integrate, then multiply (1.1) by $\frac{B_n}{\sqrt{A_n^2 + B_n^2}} \sin nt$ and integrate, and add the results:

(2.2)
$$\sqrt{A_n^2 + B_n^2} = \int_0^{2\pi} F(x(t))' \cos(nt - \delta) dt + \int_0^{2\pi} g(x(t)) \cos(nt - \delta) dt,$$

in view of (2.1). Using that x(t) is a 2π periodic solution, and Lemma 2.2, obtain

$$\int_0^{2\pi} F(x(t))' \cos(nt-\delta) dt = n \int_0^{2\pi} F(x(t)) \sin(nt-\delta) dt$$

= $n \int_{P_s} F(x(t)) \sin(nt-\delta) dt + n \int_{N_s} F(x(t)) \sin(nt-\delta) dt$
< $2n \left(F(\infty) - F(-\infty)\right)$.

Similarly, using Lemma 2.1

$$\int_0^{2\pi} g(x(t)) \cos(nt-\delta) dt < g(\infty) \int_{P_c} \cos(nt-\delta) dt + g(-\infty) \int_{N_c} \cos(nt-\delta) dt$$
$$= 2 \left(g(\infty) - g(-\infty) \right) \,.$$

The condition (1.6) follows.

Proof of Proposition 2. Consider the equation

(2.3)
$$x'' + f(x)x' + g(x) + n^2 x = E \cos nt$$

with a parameter *E*. Calculate $A_n = \int_0^{2\pi} E \cos^2 nt \, dt = E\pi$, $B_n = \int_0^{2\pi} E \cos nt \sin nt \, dt = 0$, and $\sqrt{A_n^2 + B_n^2} = E\pi$. Choose *E* so that

 \diamond

(2.4)
$$E\pi = \sqrt{A_n^2 + B_n^2} = 2n \left(F(\infty) - F(-\infty)\right) + 2 \left(g(\infty) - g(-\infty)\right) - \epsilon,$$

with $\epsilon > 0$ small, to be specified. The condition (1.6) holds for the equation (2.3). If this condition were sufficient, we would have a 2π periodic solution of (2.3), and hence

(2.5)
$$x'' + F(x)' + n^2 x = E \cos nt - g(x) \equiv \bar{e}(t) .$$

Calculate the coefficients A_n, B_n for (2.5):

(2.6)
$$\bar{A}_n = \int_0^{2\pi} \bar{e}(t) \cos nt \, dt = E\pi - \int_0^{2\pi} g(x) \cos nt \, dt, \bar{B}_n = \int_0^{2\pi} \bar{e}(t) \sin nt \, dt = -\int_0^{2\pi} g(x) \sin nt \, dt.$$

Since (2.5) is solvable, by Theorem 1.2 we have

(2.7)
$$\sqrt{\bar{A}_n^2 + \bar{B}_n^2} < 2n \left(F(\infty) - F(-\infty) \right).$$

For any $\epsilon > 0$ we can choose an index n_0 , so that for $n \ge n_0$

(2.8)
$$|\int_0^{2\pi} g(x)\sin nt\,dt| < \epsilon,$$

as follows by well known results on oscillatory integrals, see e.g., O. Costin et al [2]. Using (2.6), (2.8), followed by (2.4), obtain

$$\begin{split} \sqrt{\bar{A}_n^2 + \bar{B}_n^2} &> |\bar{A}_n| > E\pi - \epsilon \\ &= 2n \left(F(\infty) - F(-\infty) \right) + 2 \left(g(\infty) - g(-\infty) \right) - 2\epsilon \\ &> 2n \left(F(\infty) - F(-\infty) \right), \end{split}$$

contradicting (2.7), provided we fix $\epsilon < g(\infty) - g(-\infty)$.

 \diamond

We shall prove a generalization of Proposition 3 after several preliminary results. By an obvious modification of its proof, one obtains the following generalization of Proposition 1.

Proposition 4 Assume that the functions F(x) and g(x) have finite infimums and supremums on $(-\infty, \infty)$. Then the condition

$$\sqrt{A_n^2 + B_n^2} < 2n \left(\sup F - \inf F \right) + 2 \left(\sup g - \inf g \right)$$

is necessary for the existence of 2π periodic solution of (1.1).

We shall use the following result that is included in J.M. Alonso and R. Ortega [1].

Proposition 5 ([1]) Let $G(\zeta, \eta) : \mathbb{R}^2 \to \mathbb{R}^2$ be a continuous vector function, and let $V(\zeta, \eta) : \mathbb{R}^2 \to \mathbb{R}$ be a continuous function. With $\xi \equiv (\zeta, \eta)$ consider a sequence

$$\xi_{n+1} = G\left(\xi_n\right), \quad n \ge 0\,,$$

beginning with an arbitrary vector ξ_0 . Assume that

(2.9)
$$V(G(\xi)) > V(\xi), \quad \forall \xi \in \mathbb{R}^2.$$

Then $\lim_{n\to\infty} ||\xi_n|| = \lim_{n\to\infty} \sqrt{\zeta_n^2 + \eta_n^2} = \infty.$

Proof: If the sequence $\{||\xi_n||\}$ fails to tend to infinity, then $\{\xi_n\}$ has a finite accumulation point $\xi^* \in \mathbb{R}^2$. Let $\{\xi_{n_k}\}$ be a subsequence tending to ξ^* , with $n_1 < n_2 < \cdots$. Since $V(\xi_{n+1}) = V(G(\xi_n)) > V(\xi_n)$ by (2.9), the sequence $V(\xi_n)$ is increasing. Then one has

$$V(G(\xi^*)) = \lim_{k \to \infty} V(G(\xi_{n_k})) = \lim_{k \to \infty} V(\xi_{n_k+1})$$

$$\leq \lim_{k \to \infty} V(\xi_{n_{k+1}}) = V(\xi^*),$$

contradicting (2.9). (Observing that $n_k + 1 \le n_{k+1}$.)

$$\diamond$$

The next lemma says that for solution of (1.1), $x^2(t) + {x'}^2(t)$ cannot increase too much over an interval of length 2π .

Lemma 2.3 Assume assume that $e(t) \in C(R)$ is 2π periodic, the condition (1.3) holds, and moreover assume that

(2.10)
$$f(x) \ge \alpha$$
, for some $\alpha \in R$, and all $x \in R$.

Then for any initial data (x(0), x'(0)), with $c_0 = x^2(0) + n^2 {x'}^2(0)$, there is a number $c = c(c_0)$ so that the corresponding solution of (1.1) satisfies

$$x^{2}(t) + n^{2} {x'}^{2}(t) \le c$$
, for all $t \in [0, 2\pi]$.

Proof: Consider the "energy" $E(t) = \frac{1}{2}x'^2(t) + \frac{1}{2}n^2x^2(t)$. Since

$$E'(t) = -f(x)x'^{2} - g(x)x' + e(t)x'.$$

By our conditions obtain $E'(t) \leq c_1 E(t) + c_2$, and the proof follows. \diamond

We now prove the following generalization of Proposition 3.

Proposition 6 Assume that the functions F(x) and g(x) have finite infimums and supremums on $(-\infty, \infty)$ Assume also that (2.10) holds. In case

$$\sqrt{A_n^2 + B_n^2} \ge 2n \left(\sup F - \inf F \right) + 2 \left(\sup g - \inf g \right),$$

all solutions of (1.1) satisfy $\lim_{t\to\pm\infty} \left(x^2(t) + {x'}^2(t)\right) = \infty$.

Proof: Following J.M. Alonso and R. Ortega [1], we shall use Proposition 5. Given $\xi = (\zeta, \eta) \in R^2$, denote by $x(t, \xi)$ the solution of (1.1) satisfying $x(0) = \zeta$, $x'(0) = \eta$. Define the map $R^2 \to R^2$ by $G(\xi) = (x(2\pi,\xi), x'(2\pi,\xi))$, and we shall show that the sequence of iterates

$$\xi_{n+1} = G(\xi_n), \quad n = 0, 1, 2, \dots$$

is unbounded for any ξ_0 . With δ as defined by (2.1), define the function

 $V(\xi) = \eta \cos \delta - n\zeta \sin \delta + F(\zeta) \cos \delta.$

Multiply (1.1) by $\frac{A_n}{\sqrt{A_n^2 + B_n^2}} \cos nt$ and integrate, then multiply (1.1) by $\frac{B_n}{\sqrt{A_n^2 + B_n^2}} \sin nt$ and integrate, and add the results. In view of (2.1) obtain as above

(2.11)
$$\int_0^{2\pi} (x'' + n^2 x) \cos(nt - \delta) dt + \int_0^{2\pi} F(x(t))' \cos(nt - \delta) dt + \int_0^{2\pi} g(x(t)) \cos(nt - \delta) dt = \sqrt{A_n^2 + B_n^2}.$$

Integrating by parts, we express the first term on the left as

$$[x'(t)\cos(nt-\delta)]|_{0}^{2\pi} + n [x(t)\sin(nt-\delta)]|_{0}^{2\pi}$$

= [x'(2\pi) - \eta] cos \delta - n [x(2\pi) - x(\zeta)] sin \delta ,

and the second term as

$$[F(x(2\pi)) - F(x(0))]\cos \delta + n \int_0^{2\pi} F(x)\sin(nt - \delta) dt$$

We combine the non-integral terms in (2.11) as $V(G(\xi)) - V(\xi)$. Then (2.11) gives

$$\begin{split} V(G(\xi)) - V(\xi) &= \sqrt{A_n^2 + B_n^2} - n \int_0^{2\pi} F(x(t)) \sin(nt - \delta) \, dt \\ &- \int_0^{2\pi} g(x(t)) \cos(nt - \delta) \, dt \\ &> \sqrt{A_n^2 + B_n^2} - 2n \left(\sup F - \inf F \right) - 2 \left(\sup g - \inf g \right) \ge 0 \, . \end{split}$$

(The first inequality is strict because the functions g and F are non-constant by (1.2) and (1.3).) Hence, the condition (2.9) holds, and Proposition 5 applies, proving the unboundness of the sequence $(x(2n\pi,\xi), x'(2n\pi,\xi))$. If there was a sequence $\{t_k\} \to \infty$ with bounded $x'^2(t_k) + x^2(t_k)$, we would obtain a contradiction with Lemma 2.3.

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