

19 - REVIEW OF RADIATION MECHANISMS

Atomic processes can be:

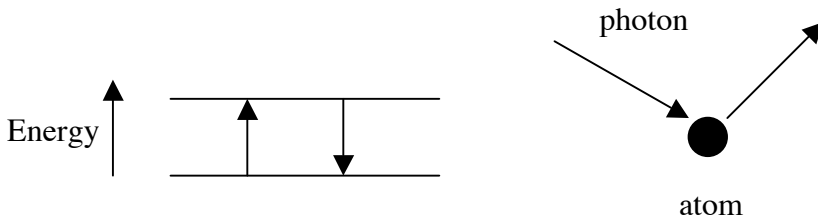
1. Scattering
2. Absorption/Thermal Emission

They may also be:

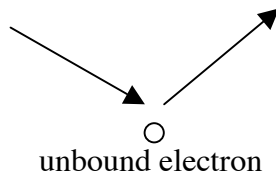
- A. Line processes
- B. Continuum processes

Scattering

Line (bound-bound) Scattering



Continuum Scattering



During the scattering process, the frequencies may change a little:

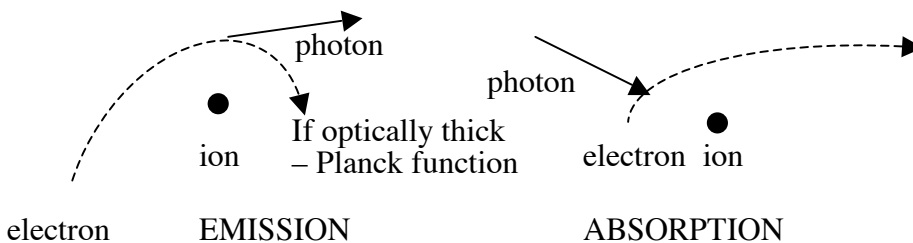
Lines \Rightarrow widths of energy levels

Continuum \Rightarrow Thomson scattering (recoil of electron) MORE BELOW!

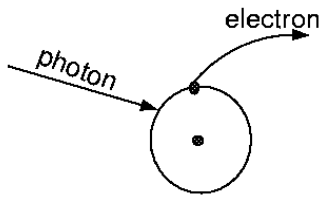
Generally, we will assume isotropic scattering, although the results won't change much when this condition is relaxed.

Absorption/Thermal Emission

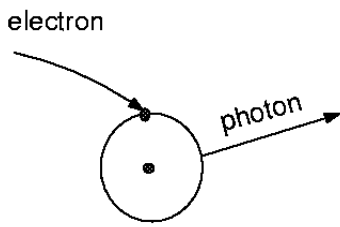
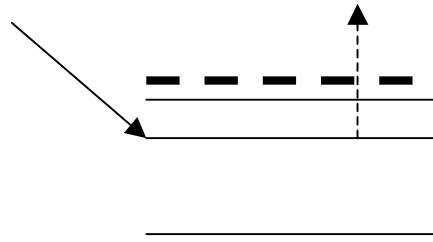
Free-free (continuum) ("Bremsstrahlung") Emission/Absorption



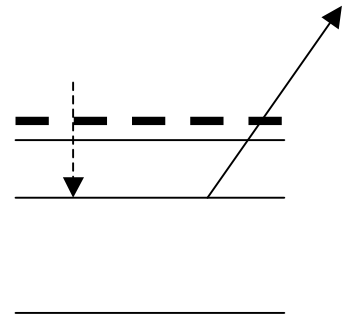
Bound-Free Processes



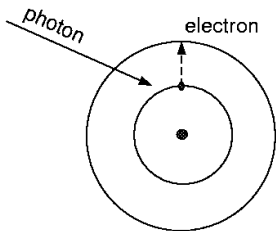
Ionization



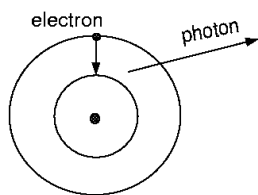
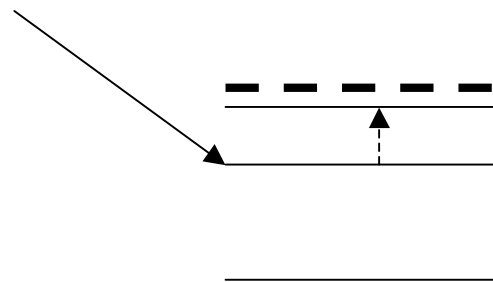
Recombination



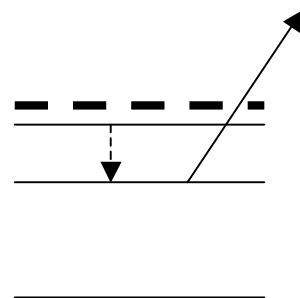
Bound-Bound Processes



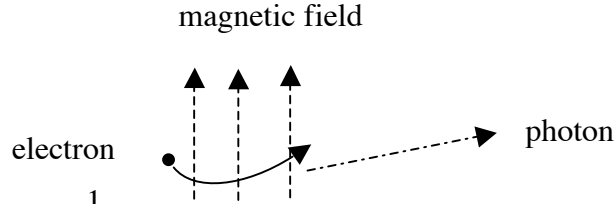
Excitation



Deexcitation



Charged Particle in a Magnetic Field



$$\text{Lorentz factor } \gamma = \frac{1}{\sqrt{1 - \left(\frac{v}{c}\right)^2}}$$

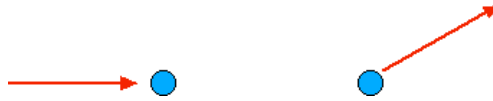
For $\gamma=1$ Cyclotron Radiation – circularly polarized

For $\gamma \gg 1$ Synchrotron Radiation – linearly polarized with $\vec{E}(\text{photons}) \perp \vec{B}(\text{field})$

For a power-law particle energy distribution, $F_\nu \propto \nu^{-\alpha}$

Scattering Regimes:

Thomson



$$\sigma_T = \frac{8\pi}{3} \left(\frac{e^2}{mc^2} \right)^2 \approx 6.65 \times 10^{-25} \text{ cm}^2 \quad \text{for } \nu < 10^{20} \text{ Hz}$$

(Klein-Nishina)

$$\sigma_{KN} \approx \frac{3}{8} \sigma_T \left(\frac{mc^2}{h\nu} \right) \left[\ln \left(\frac{2h\nu}{mc^2} \right) + \frac{1}{2} \right] \approx \frac{3 \times 10^{-5}}{\nu} \left[\ln(1.6 \times 10^{-20} \nu) + \frac{1}{2} \right] \quad \text{for } \nu \gg 10^{20} \text{ Hz}$$

Compton



$$\sigma_C \approx \begin{cases} \sigma_T & \text{for } \gamma h\nu \ll mc^2 \\ \frac{8}{3} \sigma_T \frac{mc^2}{\gamma h\nu} \left[\ln \left(\frac{2\gamma h\nu}{mc^2} \right) + \frac{1}{2} \right] & \text{for } \gamma h\nu \gg mc^2 \end{cases}$$

$$\nu_{scatt} = \nu_0 \left[1 + \left(\frac{h\nu}{mc^2} \right) (1 - \cos \phi) \right]^{-1} \quad \text{for } \gamma \ll 1 \text{ and } \gamma h\nu \ll mc^2$$

$$= \frac{4\gamma^2 \nu}{3} \quad \text{for } \gamma \gg 1 \text{ and } \gamma h\nu \ll mc^2$$

$$= \frac{\gamma mc^2}{h} \quad \text{for } \gamma \gg 1 \text{ and } \gamma h\nu \gg mc^2$$

Inverse-Compton



$$\sigma_c \approx \begin{cases} \sigma_T \left(1 - \frac{2\gamma h\nu}{mc^2} \right) & \text{and } \nu_{scatt} \approx \gamma^2 \nu_0 \quad \text{for } \gamma h\nu \ll mc^2 \\ \frac{8}{3} \sigma_T \frac{mc^2}{\gamma h\nu} \left[\ln \left(\frac{2\gamma h\nu}{mc^2} \right) + \frac{1}{2} \right] & \text{and } \nu_{scatt} \approx \frac{\gamma mc^2}{h} \quad \text{for } \gamma h\nu \gg mc^2 \end{cases}$$

Note: The inverse –Compton scattering can take place in a thermal gas or a non-thermal gas. If a synchrotron-emitting plasma has a high enough density of photons and electrons, then the synchrotron-emitting photons can inverse Compton scatter other photons nearby, This is called synchrotron self-Compton scattering (SSC). (An analogous process to free-free absorption, synchrotron self-absorption (SSA) can also occur).