

Non-existence of solutions for non-autonomous elliptic systems

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Abstract

We extend the classical Pohozaev's identity to semilinear elliptic systems of Hamiltonian type, providing a simpler approach, and a generalization, of the results of E. Mitidieri [6], R.C.A.M. Van der Vorst [14], and Y. Bozhkov and E. Mitidieri [1].

Key words: Pohozaev's identity, non-existence of solutions.

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1 Introduction

Any solution $u(x)$ of semilinear Dirichlet problem on a bounded domain $\Omega \subset \mathbb{R}^n$

$$(1.1) \quad \Delta u + f(x, u) = 0 \text{ in } \Omega, \quad u = 0 \text{ on } \partial\Omega$$

satisfies the well known Pohozaev's identity

$$(1.2) \quad \int_{\Omega} [2nF(x, u) + (2 - n)uf(x, u) + 2\sum_{i=1}^n x_i F_{x_i}(x, u)] dx = \int_{\partial\Omega} (x \cdot \nu) |\nabla u|^2 dS.$$

Here $F(x, u) = \int_0^u f(x, t) dt$, and ν is the unit normal vector on $\partial\Omega$, pointing outside. (From the equation (1.1), $\int_{\Omega} uf(x, u) dx = \int_{\Omega} |\nabla u|^2 dx$, which gives an alternative form of the Pohozaev's identity.) Pohozaev's identity is usually written for the case $f = f(u)$, but the present version is also known, see e.g., K. Schmitt [13]. A standard use of this identity is to conclude that if Ω is a star-shaped domain with respect to the origin, i.e., $x \cdot \nu \geq 0$ for all $x \in \partial\Omega$, and $f(u) = u|u|^{p-1}$, for some constant p , then the problem (1.1) has no non-trivial solutions in the super-critical case, when $p > \frac{n+2}{n-2}$. In this note we present a proof of Pohozaev's identity, which appears a little more straightforward than the usual one, see e.g., L. Evans [2], and then use a similar idea for systems, generalizing the well-known results of E. Mitidieri [6], see also R.C.A.M. Van der Vorst [14], and of Y. Bozhkov and E. Mitidieri [1], by allowing explicit dependence on x in the Hamiltonian function.

Let $z = x \cdot \nabla u = \sum_{i=1}^n x_i u_{x_i}$. It is straightforward to verify that z satisfies

$$(1.3) \quad \Delta z + f_u(x, u)z = -2f(x, u) - \sum_{i=1}^n x_i f_{x_i}(x, u).$$

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We multiply the equation (1.1) by z , and subtract from that the equation (1.3) multiplied by u , obtaining

$$(1.4) \quad \sum_{i=1}^n (zu_{x_i} - uz_{x_i})_{x_i} + \sum_{i=1}^n (f(x, u) - uf_u(x, u)) x_i u_{x_i} = 2f(x, u)u + \sum_{i=1}^n x_i f_{x_i}(x, u)u.$$

We have

$$\begin{aligned} \sum_{i=1}^n (f(x, u) - uf_u(x, u)) x_i u_{x_i} &= \sum_{i=1}^n x_i \frac{\partial}{\partial x_i} (2F - uf) - 2 \sum_{i=1}^n x_i F_{x_i} + \sum_{i=1}^n x_i f_{x_i}(x, u)u = \\ &= \sum_{i=1}^n \frac{\partial}{\partial x_i} [x_i (2F - uf)] - n(2F - uf) - 2 \sum_{i=1}^n x_i F_{x_i} + \sum_{i=1}^n x_i f_{x_i}(x, u)u. \end{aligned}$$

We then rewrite (1.4)

$$(1.5) \quad \sum_{i=1}^n [(zu_{x_i} - uz_{x_i}) + x_i (2F(x, u) - uf(x, u))]_{x_i} = 2nF(x, u) + (2-n)uf(x, u) + 2 \sum_{i=1}^n x_i F_{x_i}.$$

Integrating over Ω , we conclude the Pohozaev's identity (1.2). (The only non-zero boundary term is $\sum_{i=1}^n \int_{\partial\Omega} zu_{x_i} \nu_i dS$. Since $\partial\Omega$ is a level set of u , $\nu = \pm \frac{\nabla u}{|\nabla u|}$, i.e., $u_{x_i} = \pm |\nabla u| \nu_i$. Then $z = \pm (x \cdot \nu) |\nabla u|$, and $\sum_{i=1}^n u_{x_i} \nu_i = \pm |\nabla u|$.)

We refer to (1.5) as a *differential form* of Pohozaev's identity. For radial solutions on a ball, the corresponding version of (1.5) played a crucial role in the study of exact multiplicity of solutions, see T. Ouyang and J. Shi [7], and also P. Korman [5], which shows the potential usefulness of this identity.

2 Non-existence of solutions for a class of systems

The following class of systems has attracted considerable attention recently

$$(2.1) \quad \begin{aligned} \Delta u + H_v(u, v) &= 0 \text{ in } \Omega, \quad u = 0 \text{ on } \partial\Omega \\ \Delta v + H_u(u, v) &= 0 \text{ in } \Omega, \quad v = 0 \text{ on } \partial\Omega, \end{aligned}$$

where $H(u, v)$ is a given differentiable function, see e.g., the following surveys: D.G. de Figueiredo [3], P. Quittner and P. Souplet [11], B. Ruf [12], see also P. Korman [4]. This system is of *Hamiltonian* type, so that it has some of the properties of scalar equations.

More generally, let $H = H(x, u_1, u_2, \dots, u_m, v_1, v_2, \dots, v_m)$, with integer $m \geq 1$, and consider the Hamiltonian system of $2m$ equations

$$(2.2) \quad \begin{aligned} \Delta u_k + H_{v_k} &= 0 \text{ in } \Omega, \quad u_k = 0 \text{ on } \partial\Omega, \quad k = 1, 2, \dots, m \\ \Delta v_k + H_{u_k} &= 0 \text{ in } \Omega, \quad v_k = 0 \text{ on } \partial\Omega, \quad k = 1, 2, \dots, m. \end{aligned}$$

We call solution of (2.2) to be positive, if $u_k(x) > 0$ and $v_k(x) > 0$ for all $x \in \Omega$, and all k . We consider only the classical solutions, with u_k and v_k of class $C^2(\Omega) \cap C^1(\bar{\Omega})$. We have the following generalization of the results of [1] and [6].

Theorem 2.1 *Assume that $H(x, u_1, u_2, \dots, u_m, v_1, v_2, \dots, v_m) \in C^2(\Omega \times R_+^m \times R_+^m) \cap C(\bar{\Omega} \times \bar{R}_+^m \times \bar{R}_+^m)$ satisfies*

$$(2.3) \quad H(x, 0, \dots, 0, 0, \dots, 0) = 0 \text{ for all } x \in \partial\Omega.$$

Then for any positive solution of (2.2), and any real numbers a_1, \dots, a_m , one has

$$(2.4) \quad \begin{aligned} \int_{\Omega} [2nH + (2-n) \sum_{k=1}^m (a_k u_k H_{u_k} + (2-a_k) v_k H_{v_k}) + 2 \sum_{i=1}^n x_i H_{x_i}] dx \\ = 2 \sum_{k=1}^m \int_{\partial\Omega} (x \cdot \nu) |\nabla u_k| |\nabla v_k| dS. \end{aligned}$$

Proof: Define $p_k = x \cdot \nabla u_k = \sum_{i=1}^n x_i u_{kx_i}$, and $q_k = x \cdot \nabla v = \sum_{i=1}^n x_i v_{kx_i}$, $k = 1, 2, \dots, m$. These functions satisfy the system

$$(2.5) \quad \begin{aligned} \Delta p_k + \sum_{j=1}^m H_{v_k u_j} p_j + \sum_{j=1}^m H_{v_k v_j} q_j &= -2H_{v_k} - \sum_{i=1}^n x_i H_{v_k x_i}, \quad k = 1, 2, \dots, m \\ \Delta q_k + \sum_{j=1}^m H_{u_k u_j} p_j + \sum_{j=1}^m H_{u_k v_j} q_j &= -2H_{u_k} - \sum_{i=1}^n x_i H_{u_k x_i}, \quad k = 1, 2, \dots, m. \end{aligned}$$

We multiply the first equation in (2.2) by q_k , and subtract from that the first equation in (2.5) multiplied by v_k . The result can be written as

$$(2.6) \quad \begin{aligned} &\sum_{i=1}^n [(u_{kx_i} q_k - p_{kx_i} v_k)_{x_i} + (-u_{kx_i} q_{kx_i} + v_{kx_i} p_{kx_i})] \\ &+ H_{v_k} q_k - \sum_{j=1}^m H_{v_k u_j} p_j v_k - \sum_{j=1}^m H_{v_k v_j} q_j v_k = 2v_k H_{v_k} + v_k \sum_{i=1}^n x_i H_{v_k x_i}. \end{aligned}$$

Similarly, we multiply the second equation in (2.2) by p_k , and subtract from that the second equation in (2.5) multiplied by u_k , and write the result as

$$(2.7) \quad \begin{aligned} &\sum_{i=1}^n [(v_{kx_i} p_k - q_{kx_i} u_k)_{x_i} + (-v_{kx_i} p_{kx_i} + u_{kx_i} q_{kx_i})] \\ &+ H_{u_k} p_k - \sum_{j=1}^m H_{u_k u_j} p_j u_k - \sum_{j=1}^m H_{u_k v_j} q_j u_k = 2u_k H_{u_k} + u_k \sum_{i=1}^n x_i H_{u_k x_i}. \end{aligned}$$

Adding the equations (2.6) and (2.7), we get

$$\begin{aligned} &\sum_{i=1}^n [u_{kx_i} q_k - p_{kx_i} v_k + v_{kx_i} p_k - q_{kx_i} u_k]_{x_i} + H_{u_k} p_k + H_{v_k} q_k - \sum_{j=1}^m H_{u_k u_j} p_j u_k \\ &- \sum_{j=1}^m H_{u_k v_j} q_j u_k - \sum_{j=1}^m H_{v_k u_j} p_j v_k - \sum_{j=1}^m H_{v_k v_j} q_j v_k \\ &= 2u_k H_{u_k} + 2v_k H_{v_k} + u_k \sum_{i=1}^n x_i H_{u_k x_i} + v_k \sum_{i=1}^n x_i H_{v_k x_i}. \end{aligned}$$

We now sum in k , putting the result into the form

$$\begin{aligned} &\sum_{k=1}^m \sum_{i=1}^n [u_{kx_i} q_k - p_{kx_i} v_k + v_{kx_i} p_k - q_{kx_i} u_k]_{x_i} \\ &+ \sum_{i=1}^n x_i (2H - \sum_{k=1}^m u_k H_{u_k} - \sum_{k=1}^m v_k H_{v_k})_{x_i} = 2 \sum_{k=1}^m u_k H_{u_k} + 2 \sum_{k=1}^m v_k H_{v_k} + 2 \sum_{i=1}^n x_i H_{x_i}. \end{aligned}$$

Writing,

$$\begin{aligned} \sum_{i=1}^n x_i \frac{\partial}{\partial x_i} (2H - \sum_{k=1}^m u_k H_{u_k} - \sum_{k=1}^m v_k H_{v_k}) &= \sum_{i=1}^n \frac{\partial}{\partial x_i} [x_i (2H - \sum_{k=1}^m u_k H_{u_k} - \sum_{k=1}^m v_k H_{v_k})] \\ &- n(2H - \sum_{k=1}^m u_k H_{u_k} - \sum_{k=1}^m v_k H_{v_k}), \end{aligned}$$

we obtain the differential form of Pohozaev's identity

$$\begin{aligned} &\sum_{k=1}^m \sum_{i=1}^n [u_{kx_i} q_k - p_{kx_i} v_k + v_{kx_i} p_k - q_{kx_i} u_k + x_i (2H - \sum_{k=1}^m u_k H_{u_k} - \sum_{k=1}^m v_k H_{v_k})]_{x_i} \\ &= 2nH + (2 - n) (\sum_{k=1}^m u_k H_{u_k} + \sum_{k=1}^m v_k H_{v_k}) + 2 \sum_{i=1}^n x_i H_{x_i}. \end{aligned}$$

Integrating, we obtain, in view of (2.3),

$$(2.8) \quad \begin{aligned} &\int_{\Omega} [2nH(u, v) + (2 - n) (\sum_{k=1}^m u_k H_{u_k} + \sum_{k=1}^m v_k H_{v_k}) + 2 \sum_{i=1}^n x_i H_{x_i}] dx \\ &= 2 \sum_{k=1}^m \int_{\partial\Omega} (x \cdot \nu) |\nabla u_k| |\nabla v_k| dS. \end{aligned}$$

(Since we consider positive solutions, and $\partial\Omega$ is a level set for both u_k and v_k , we have $\nu = -\frac{\nabla u_k}{|\nabla u_k|} = -\frac{\nabla v_k}{|\nabla v_k|}$, i.e., $u_{ki} = -|\nabla u_k| \nu_i$ and $v_{ki} = -|\nabla v_k| \nu_i$ on the boundary $\partial\Omega$.) From the first equation in (2.2), $\int_{\Omega} v_k H_{v_k} dx = \int_{\Omega} \nabla u_k \cdot \nabla v_k dx$, while from the second equation $\int_{\Omega} u_k H_{u_k} dx = \int_{\Omega} \nabla u_k \cdot \nabla v_k dx$, i.e., for each k

$$\int_{\Omega} v_k H_{v_k} dx = \int_{\Omega} u_k H_{u_k} dx.$$

Using this in (2.8), we conclude the proof. \diamond

Remarks

1. We consider only the classical solutions. Observe that by our conditions and elliptic regularity, classical solutions are in fact of class $C^3(\Omega)$, so that all quantities in the above proof are well defined.
2. In case H is independent of x , the condition (2.3) can be assumed without loss of generality.

As a consequence, we have the following non-existence result.

Proposition 1 *Assume that Ω is a star-shaped domain with respect to the origin, and for some real constants $\alpha_1, \dots, \alpha_m$, all $u_k > 0$, $v_k > 0$, and all $x \in \Omega$, we have*

$$(2.9) \quad nH + (2-n)\sum_{k=1}^m (\alpha_k u_k H_{u_k} + (1-\alpha_k)v_k H_{v_k}) + \sum_{i=1}^n x_i H_{x_i} < 0.$$

Then the problem (2.2) has no positive solutions.

Proof: We use the identity (2.4), with $a_k/2 = \alpha_k$. Then, assuming existence of positive solution, the left hand side of (2.4) is negative, while the right hand side is non-negative, a contradiction. \diamond

Observe, that it suffices to assume that Ω is star-shaped with respect to any one of its points (which we then take to be the origin).

In case $m = 1$, and $H = H(u, v)$, we recover the following condition of E. Mitidieri [6].

Proposition 2 *Assume that Ω is a star-shaped domain with respect to the origin, and for some real constant α , and all $u > 0$, $v > 0$ we have*

$$(2.10) \quad \alpha u H_u(u, v) + (1-\alpha)v H_v(u, v) > \frac{n}{n-2} H(u, v).$$

Then the problem (2.1) has no positive solutions.

Comparing this result to E. Mitidieri [6], observe that we do not require that $H_u(0, 0) = H_v(0, 0) = 0$.

An important subclass of (2.1) is

$$(2.11) \quad \begin{aligned} \Delta u + f(v) &= 0 \text{ in } \Omega, \quad u = 0 \text{ on } \partial\Omega \\ \Delta v + g(u) &= 0 \text{ in } \Omega, \quad v = 0 \text{ on } \partial\Omega, \end{aligned}$$

which corresponds to $H(u, v) = F(v) + G(u)$, where $F(v) = \int_0^v f(t) dt$, $G(u) = \int_0^u g(t) dt$. Unlike [6], we do not require that $f(0) = g(0) = 0$. The Theorem 2.1 now reads as follows.

Theorem 2.2 *Let $f, g \in C(\bar{R}_+)$. For any positive solution of (2.11), and any real number a , one has*

$$(2.12) \quad \begin{aligned} \int_{\Omega} [2n(F(v) + G(u)) + (2-n)(avf(v) + (2-a)ug(u))] dx \\ = 2 \int_{\partial\Omega} (x \cdot \nu) |\nabla u| |\nabla v| dS. \end{aligned}$$

More generally, we consider

$$(2.13) \quad \begin{aligned} \Delta u + f(x, v) &= 0 \text{ in } \Omega, \quad u = 0 \text{ on } \partial\Omega \\ \Delta v + g(x, u) &= 0 \text{ in } \Omega, \quad v = 0 \text{ on } \partial\Omega, \end{aligned}$$

with $H(x, u, v) = F(x, v) + G(x, u)$, where $F(x, v) = \int_0^v f(x, t) dt$, $G(x, u) = \int_0^u g(x, t) dt$.

Theorem 2.3 Let $f, g \in C(\Omega \times \bar{R}_+)$. For any positive solution of (2.13), and any real number a , one has

$$(2.14) \int_{\Omega} [2n(F(x, v) + G(x, u)) + (2 - n)(avf(x, v) + (2 - a)ug(x, u)) + 2\sum_{i=1}^n x_i (F_{x_i} + G_{x_i})] dx \\ = 2 \int_{\partial\Omega} (x \cdot \nu) |\nabla u| |\nabla v| dS.$$

We now consider a particular system

$$(2.15) \quad \begin{aligned} \Delta u + v^p &= 0 \text{ in } \Omega, \quad u = 0 \text{ on } \partial\Omega \\ \Delta v + g(x, u) &= 0 \text{ in } \Omega, \quad v = 0 \text{ on } \partial\Omega, \end{aligned}$$

with $g(x, u) \in C(\Omega \times \bar{R}_+)$, and a constant $p > 0$.

Theorem 2.4 Assume that Ω is a star-shaped domain with respect to the origin, and

$$(2.16) \quad nG(x, u) + (2 - n) \left(1 - \frac{n}{(n-2)(p+1)} \right) ug(x, u) + \sum_{i=1}^n x_i G_{x_i} < 0, \text{ for } x \in \Omega, \text{ and } u > 0.$$

Then the problem (2.15) has no positive solutions.

Proof: We use the identity (2.14), with $f(v) = v^p$. We select the constant a , so that

$$2nF(v) + (2 - n)avf(v) = 0,$$

i.e., $a = \frac{2n}{(n-2)(p+1)}$. Then, assuming existence of a positive solution, the left hand side of (2.14) is negative, while the right hand side is non-negative, a contradiction. \diamond

Observe that in case $p = 1$, the Theorem 2.4 provides a non-existence result for a biharmonic problem with Navier boundary conditions

$$(2.17) \quad \Delta^2 u = g(x, u) \text{ in } \Omega, \quad u = \Delta u = 0 \text{ on } \partial\Omega.$$

Proposition 3 Assume that Ω is a star-shaped domain with respect to the origin, and the condition (2.16), with $p = 1$, holds. Then the problem (2.17) has no positive solutions.

Finally, we consider the system

$$(2.18) \quad \begin{aligned} \Delta u + v^p &= 0 \text{ in } \Omega, \quad u = 0 \text{ on } \partial\Omega \\ \Delta v + u^q &= 0 \text{ in } \Omega, \quad v = 0 \text{ on } \partial\Omega. \end{aligned}$$

The curve $\frac{1}{p+1} + \frac{1}{q+1} = \frac{n-2}{n}$ is called a *critical hyperbola*. We recover the following well known result of E. Mitidieri [6], see also R.C.A.M. Van der Vorst [14]. (Observe that we relax the restriction $p, q > 1$ from [6].)

Proposition 4 Assume that $p, q > 0$, and

$$(2.19) \quad \frac{1}{p+1} + \frac{1}{q+1} < \frac{n-2}{n}.$$

Then the problem (2.18) has no positive solutions.

Proof: Condition (2.19) implies (2.16), and then the Theorem 2.4 applies. \diamond

In case $p = 1$, we recover the following known result, see E. Mitidieri [6].

Proposition 5 Assume that Ω is a star-shaped domain with respect to the origin, and $q > \frac{n+4}{n-4}$. Then the problem

$$(2.20) \quad \Delta^2 u = u^q \text{ in } \Omega, \quad u = \Delta u = 0 \text{ on } \partial\Omega$$

has no positive solutions.

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